# Semiconductor Devices

Introduction to Semiconductors



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  - $\,\sim\,200\,keV:$  photon energy for X-rays used in medical diagnosis

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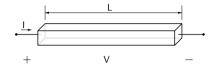
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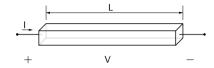
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  - 70 Hz to 150 kHz: frequency range of hearing for whales and dolphins
  - 535 kHz to 1605 kHz: AM radio frequency range (Medium Wave)

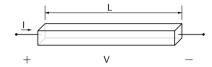
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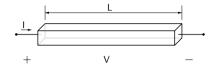




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Material	Туре	$\sigma  (\Omega ext{-cm})^{-1}$
Copper	conductor	$\sim 6  imes 10^5$
Glass	insulator	$10^{-17}$ to $10^{-13}$
Silicon	semiconductor	$\sim 10^{-5}$

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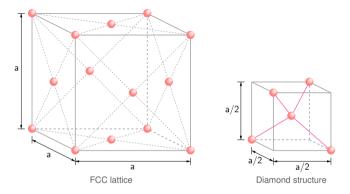


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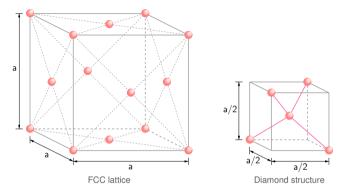
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Polycrystalline and amorphous semiconductors have relatively poor properties, but the ease of manufacturing and low cost makes them attractive for certain applications. e.g., as solar cells, thin-film transistors for display devices.

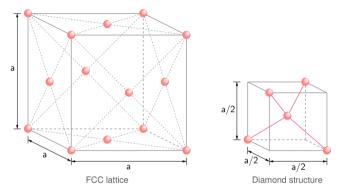




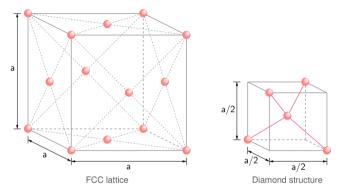
\* Si and GaAs have the diamond structure.



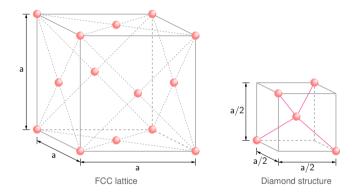
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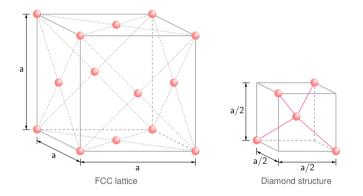


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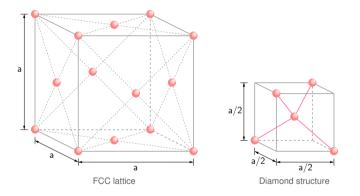


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- \* The structure of GaAs is similar. Each Ga atom has four As neighbours, and vice versa.



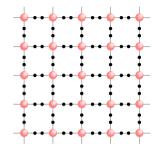


\* In silicon, the distance between neighbouring atoms is  $\sqrt{3} a/4$ , with a = 5.43 Å.



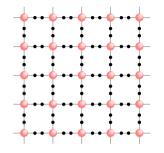
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## Electrons and holes: the bond picture



2-D representation of silicon lattice

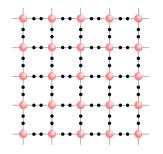
## Electrons and holes: the bond picture



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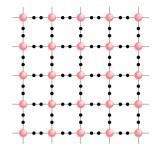
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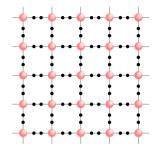
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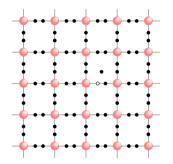
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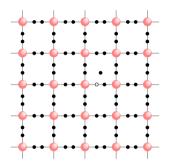
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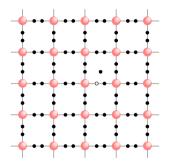
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- \* At 0 K, all valence electrons are held by the covalent bonds, no electrons are available for conduction, and the material behaves like an insulator.

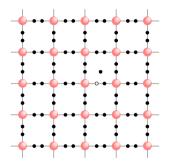




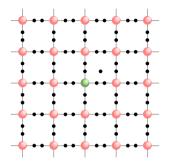
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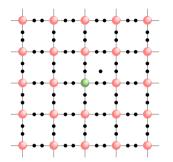


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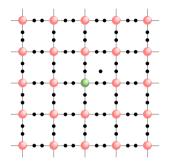


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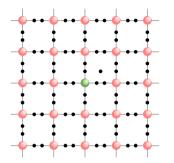




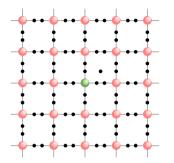
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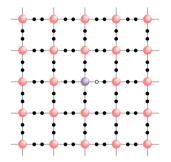
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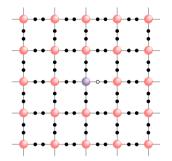


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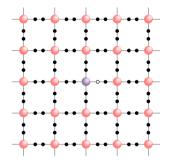


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- \* We say that the group V atom has "donated" a free electron to the lattice which is available for conduction.

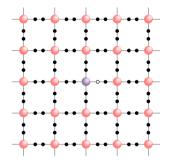




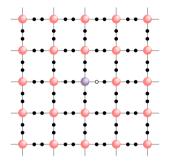
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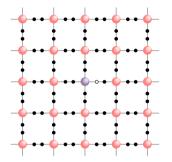
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- \* We say that the group III atom has "accepted" an electron from a Si-Si bond, which is equivalent to transferring the vacancy (hole) to that bond.

The bond picture of a semiconductor gives us some insight, but it leaves several questions unanswered.

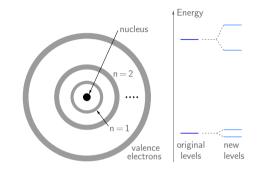
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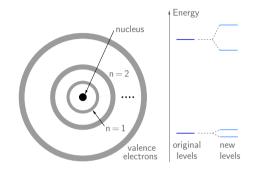
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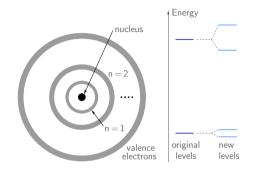
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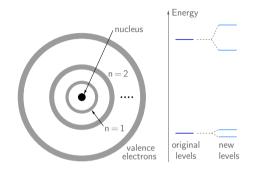




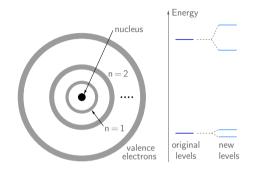
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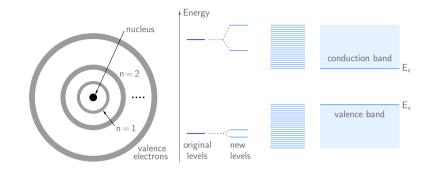
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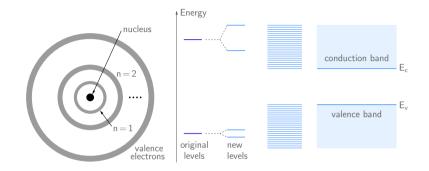


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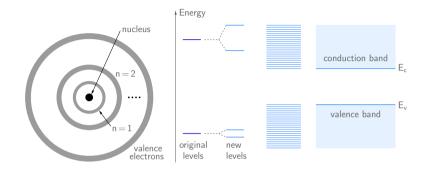


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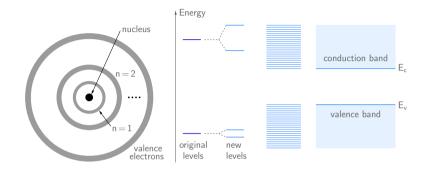




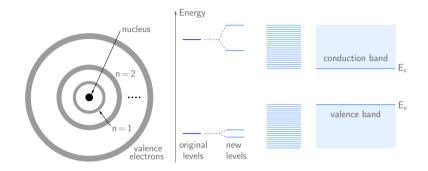
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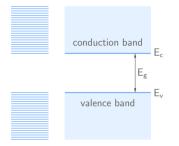
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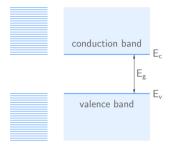


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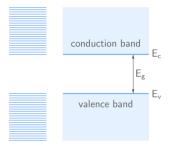


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- \* For semiconductors, the states get bunched such that, in a certain energy range,  $E_v < E < E_c$ , there are no states at all.

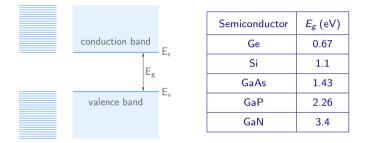




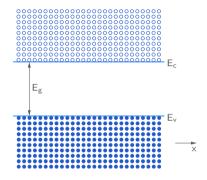
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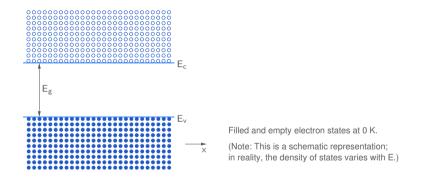


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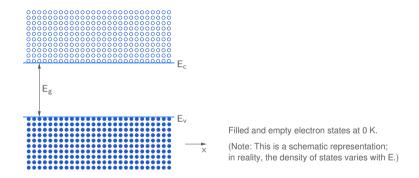


Filled and empty electron states at 0 K.

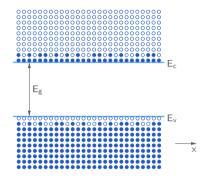
(Note: This is a schematic representation; in reality, the density of states varies with E.)



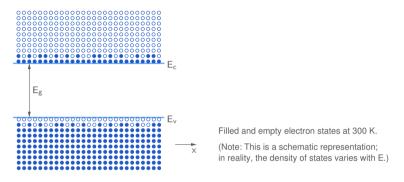
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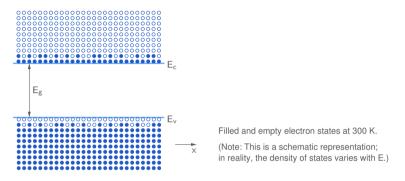
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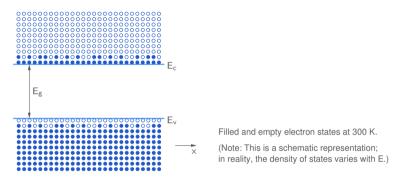
Filled and empty electron states at 300 K. (Note: This is a schematic representation; in reality, the density of states varies with E.)



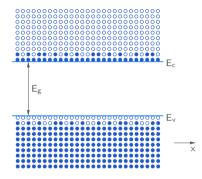
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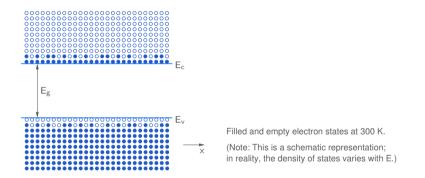


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- \* The density of electrons (or holes) in the above situation is denoted by  $n_i$ , the "intrinsic carrier concentration," and it is about  $10^{10} \text{ cm}^{-3}$  for Si at T = 300 K. (Note that it is much smaller than the density of silicon atoms in the crystal, i.e.,  $5 \times 10^{22} \text{ cm}^{-3}$ ).

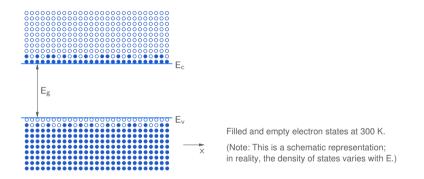


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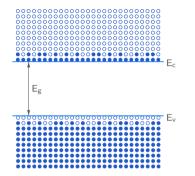
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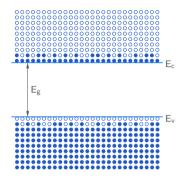


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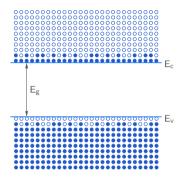


- \* An electron in the conduction band can move to one of the large number of empty states in the conduction band and contribute to a current.
- \* Similarly, a hole in the valence band can move to one of the large number of filled states in the valence band and contribute to a current.

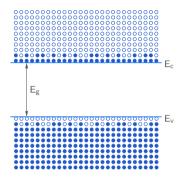




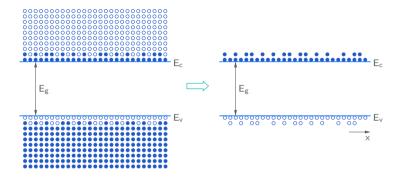
\* Filled states in the conduction band (the mobile electrons or simply "electrons") and empty states in the valence band (the mobile vacancies or "holes") are confined to a narrow energy range near  $E_c$  and  $E_v$ , respectively ( $\sim 100 \text{ meV}$  at 300 K).



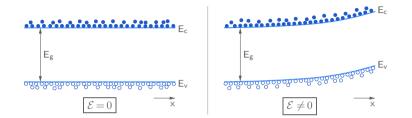
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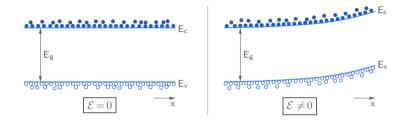


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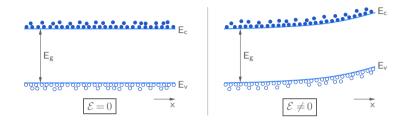


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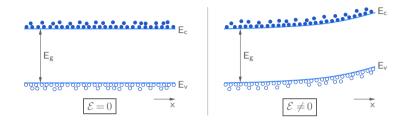




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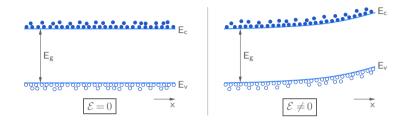


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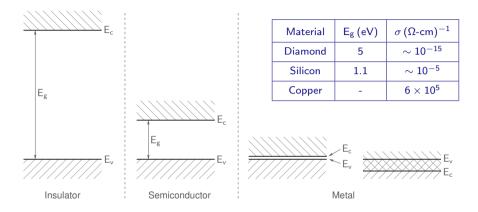


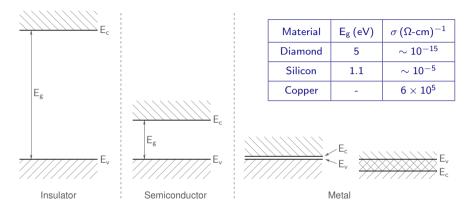
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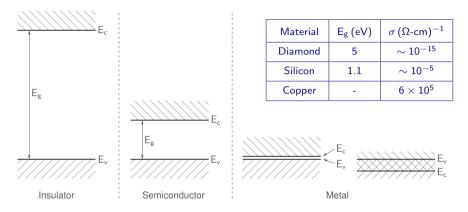


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- \* The constant in the above equation is irrelevant because only differences such as  $(E_c(x_1) E_c(x_2))$  are important, and the constant drops out.

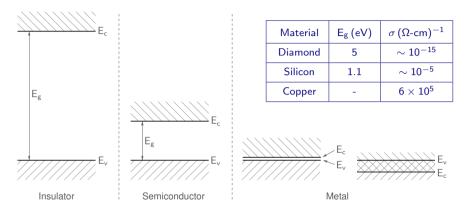




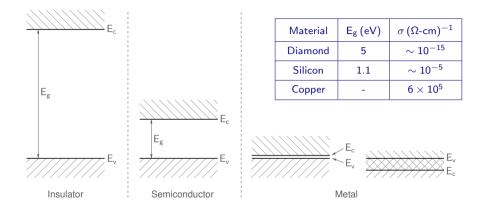
\* The electrical conductivity of a crystalline material depends on its energy gap  $E_g$ .

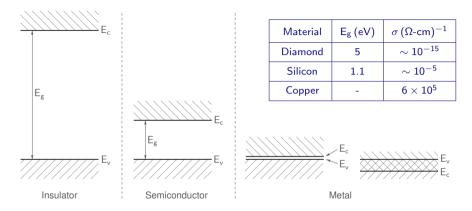


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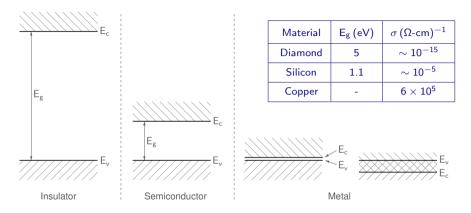


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- \* In an insulator,  $E_g$  is so large that there are no electrons (in the conduction band) or holes (in the valence band) at room temperature  $\rightarrow$  low conductivity.
- \* In a metal,  $E_g$  is either very small or non-existent. As a result, electrons in the filled states can move to one of the large number of vacant states  $\rightarrow$  high conductivity.

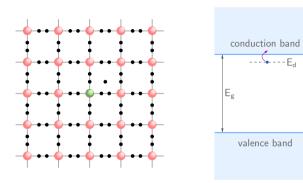




\* For a semiconductor, the situation is between these two extremes  $\rightarrow$  moderate conductivity.

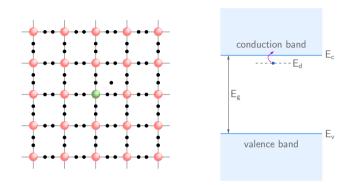


- \* For a semiconductor, the situation is between these two extremes  $\rightarrow$  moderate conductivity.
- \* Note that we have only looked at an "intrinsic" semiconductor. With addition of appropriate impurity (donor or acceptor) atoms, the conductivity of a semiconductor can be changed very significantly.

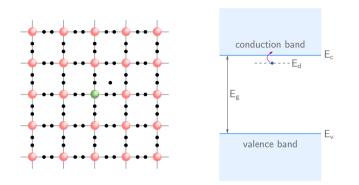


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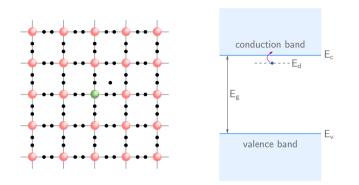
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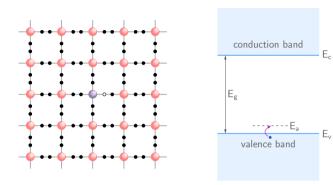
\* When a silicon atom is replaced with a donor (group V) atom, it introduces an energy level (state) in the forbidden gap, which is close to  $E_c$ .

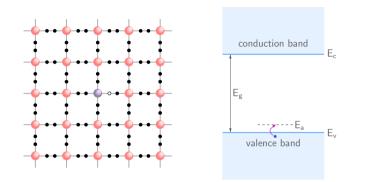


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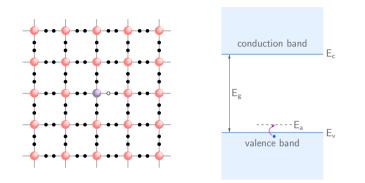


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- \* At high temperatures, the electron can cross the energy barrier  $(E_c E_d)$  and enter the conduction band.

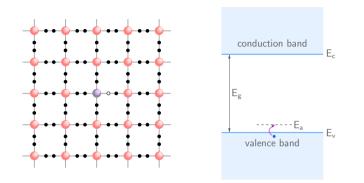




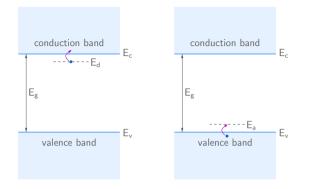
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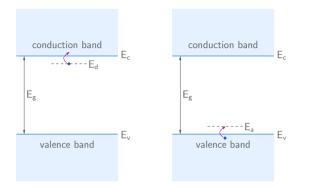


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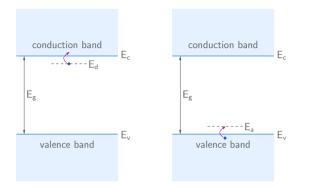


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- \* At low temperatures, the acceptor state is empty, i.e., there is a vacancy around the acceptor atom.
- \* At high temperatures, an electron from the valence band can cross the energy barrier  $(E_a E_v)$ , leaving a hole in the valence band.

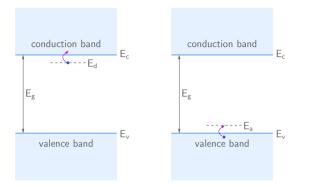




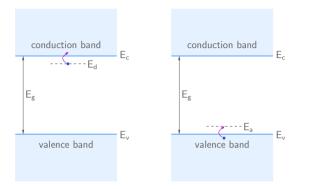
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