

SEMICONDUCTOR DEVICES

p-n Junctions: Part 1



M. B. Patil

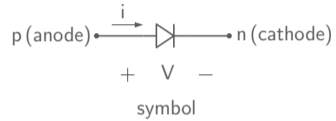
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Indian Institute of Technology Bombay

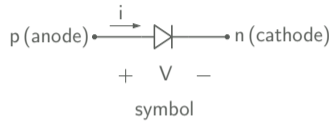


schematic diagram



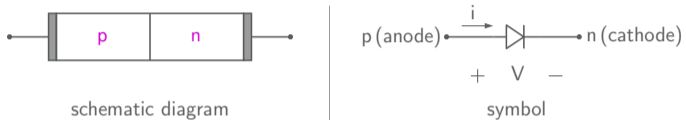


schematic diagram

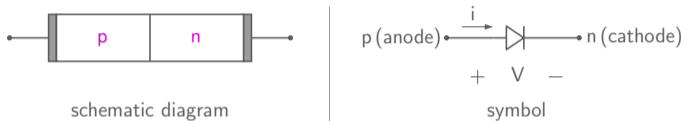


symbol

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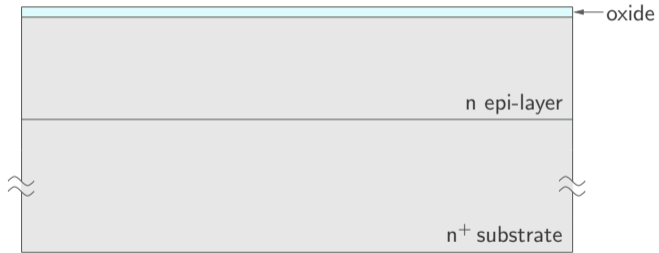
- * A p - n junction is useful as a stand-alone device (the diode).
- * It is also an integral part of devices such as transistors, IGBTs, thyristors, etc.
- * In integrated circuits, pn junctions are used to provide isolation between devices.
- * We will focus on semiconductor p - n junctions first and look at metal-semiconductor junctions later.



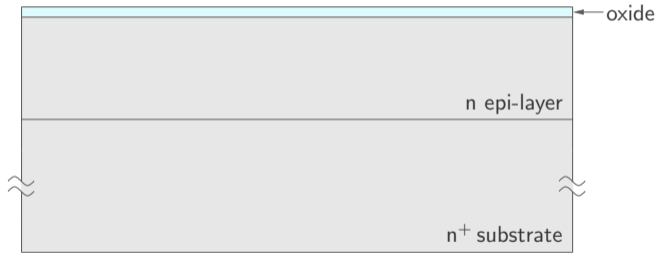
Start with n^+ substrate, with n epitaxial layer grown on top.



Deposit SiO₂.



Deposit SiO_2 .



Apply photoresist.



Apply photoresist.



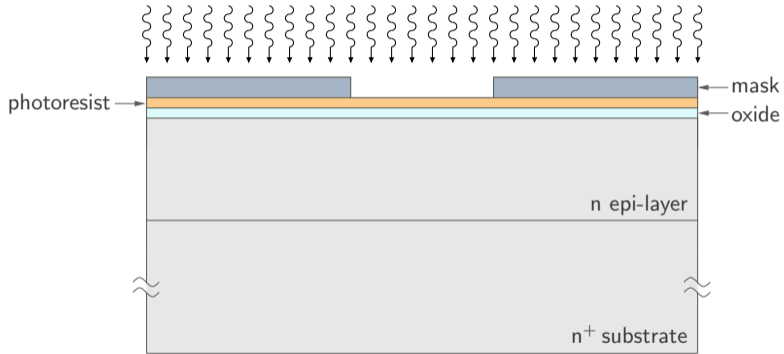
Place mask.



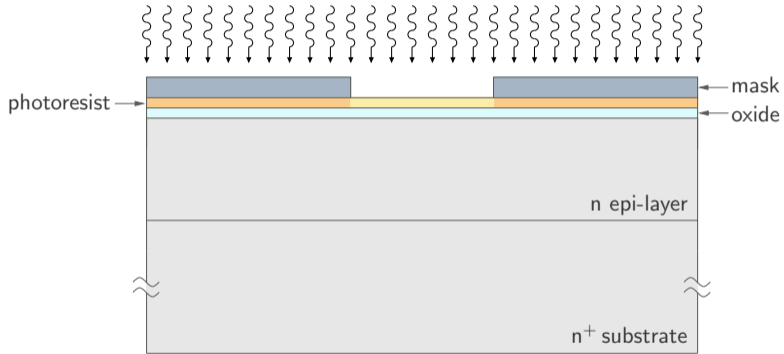
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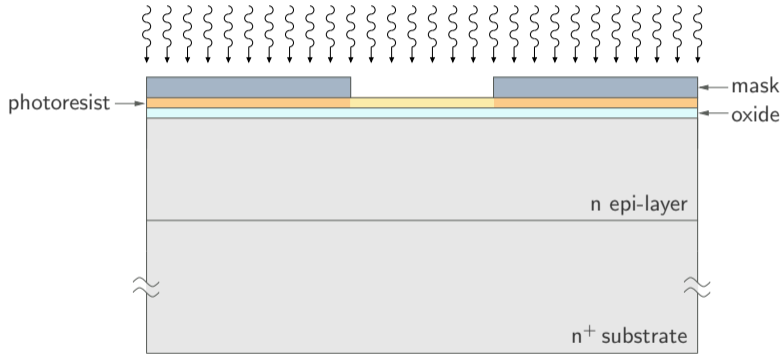
Expose to UV light.



Expose to UV light.



Expose to UV light.



Remove mask.



Remove mask.



Develop photoresist.



Develop photoresist.



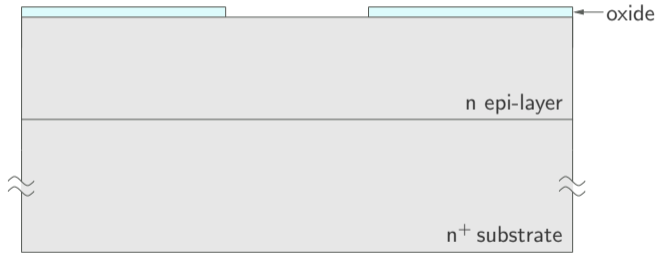
Etch oxide (in HF).



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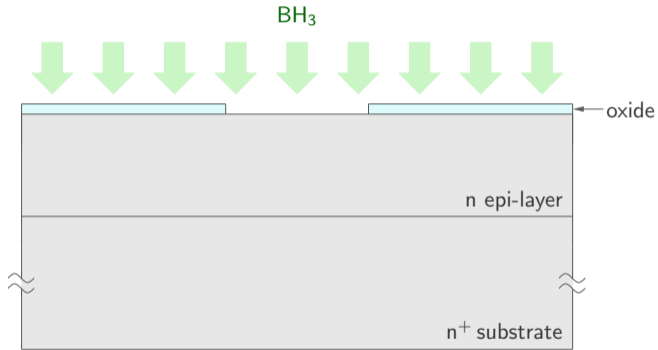
Remove photoresist.



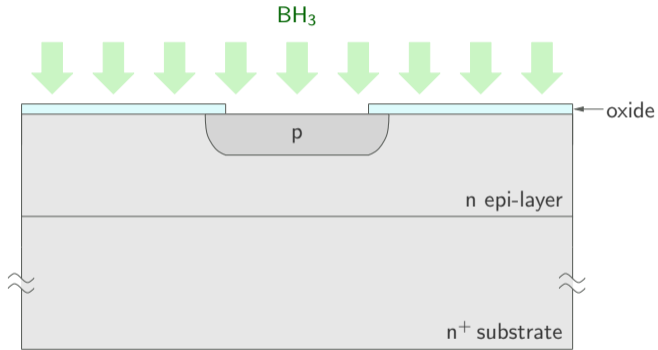
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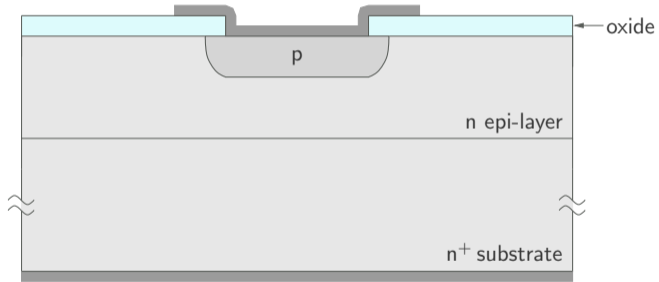
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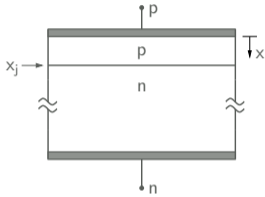


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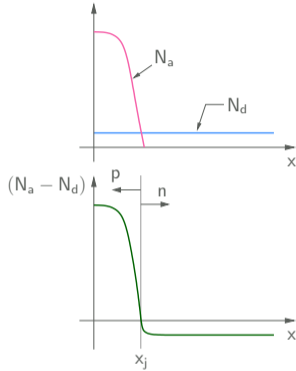


Add metal contacts (a few steps).

Idealised p - n junction diode structure

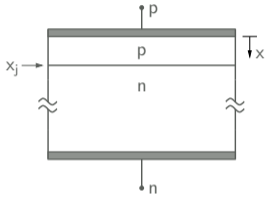


Fabricated structure

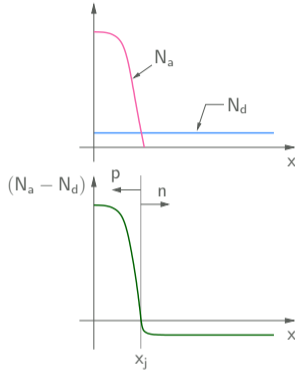


Doping densities

Idealised p - n junction diode structure



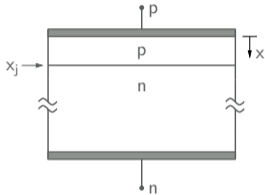
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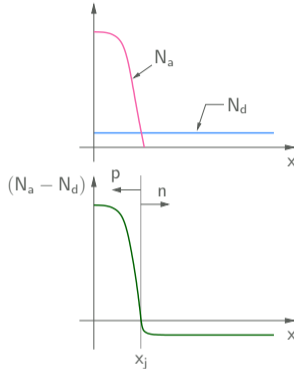
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- * For our analysis, we will consider a simplified structure with p -type doping on one side and n -type on the other.

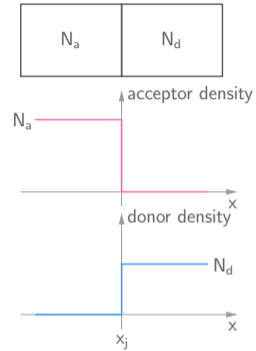
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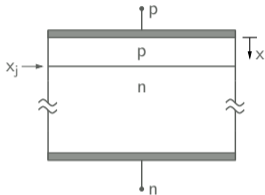
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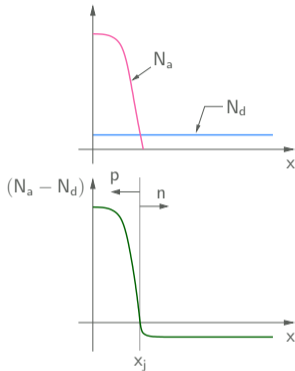
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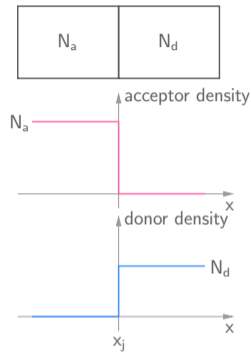
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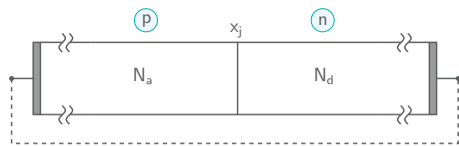
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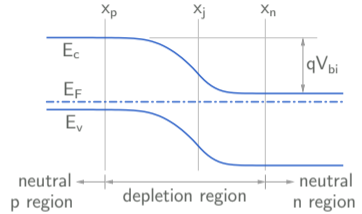
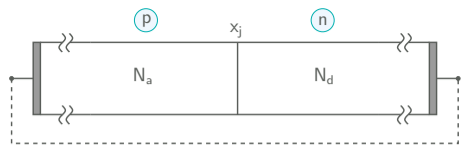
Idealised structure

- * For our analysis, we will consider a simplified structure with p -type doping on one side and n -type on the other.
- * We will assume the doping densities to change abruptly at the junction \rightarrow "abrupt" pn junction.

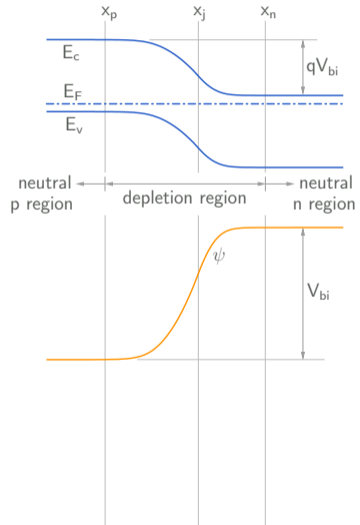
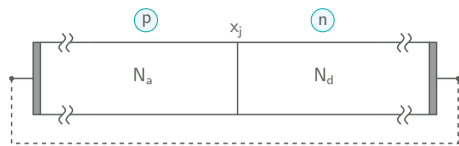
pn junction in equilibrium



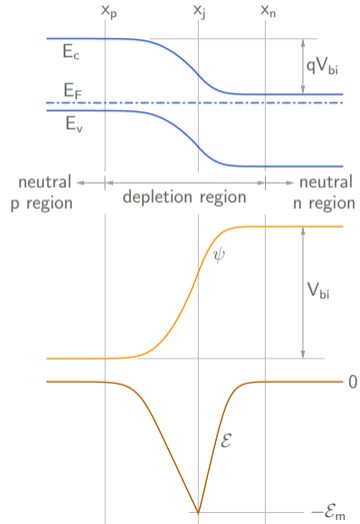
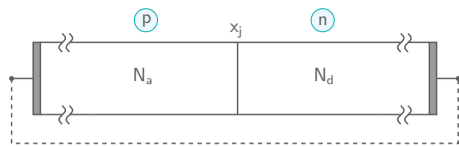
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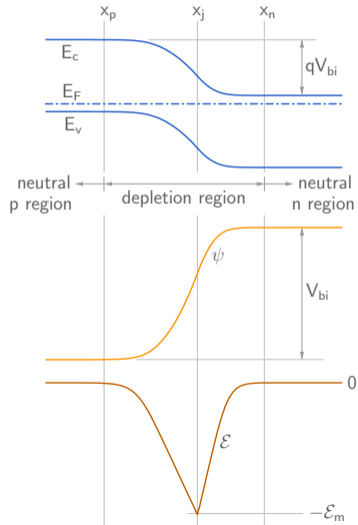


pn junction in equilibrium



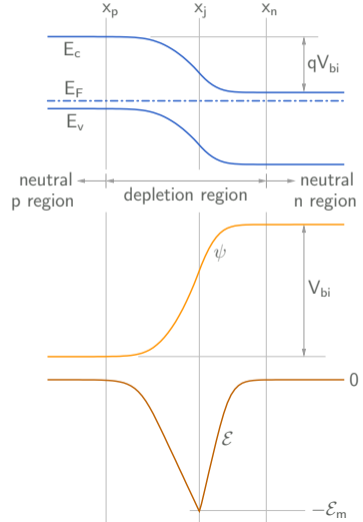
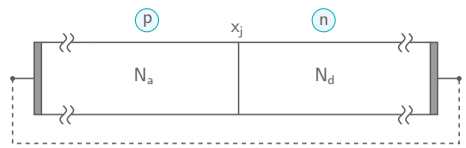
pn junction in equilibrium

* There is a “depletion region” in which the potential ψ varies.



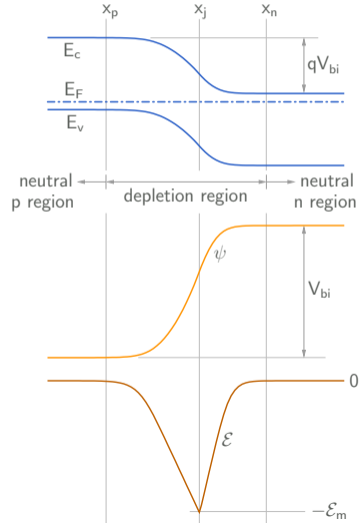
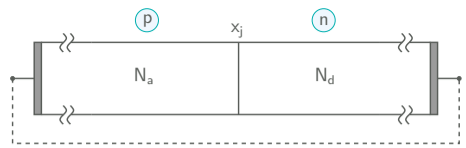
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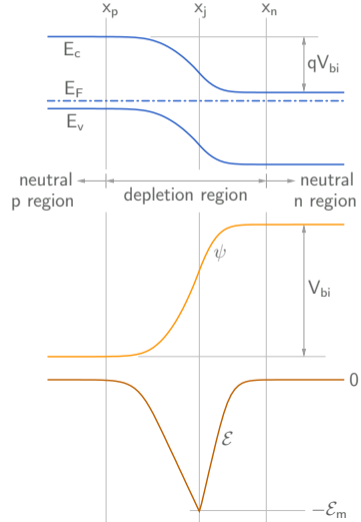
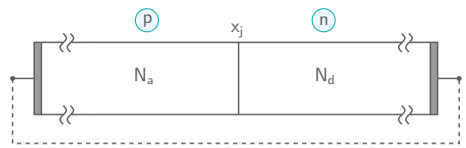
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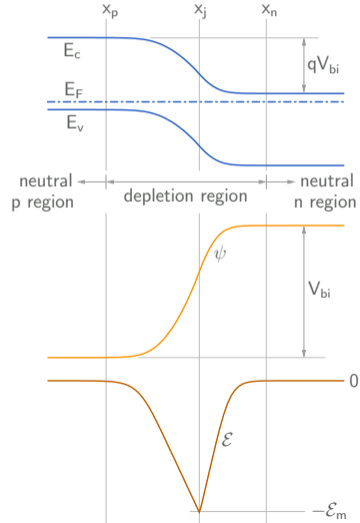
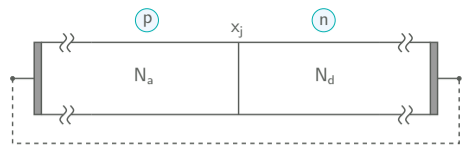
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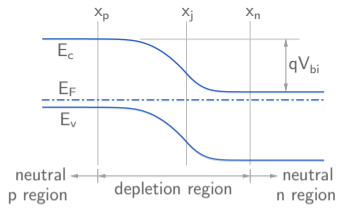


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- * Let us check if this picture is consistent with Poisson's equation.



pn junction in equilibrium

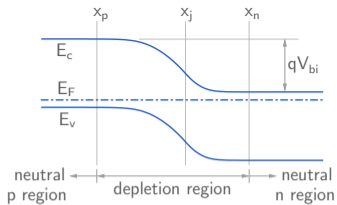


pn junction in equilibrium

Charge density:

$$p(x) = N_v \exp - \left(\frac{E_F - E_v(x)}{kT} \right),$$

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pn junction in equilibrium

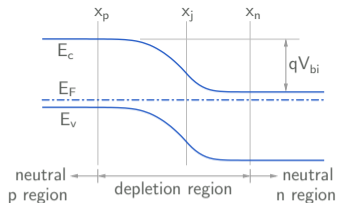
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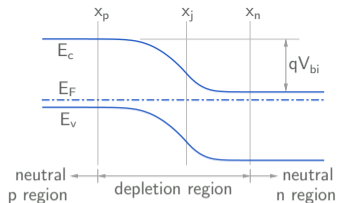


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pn junction in equilibrium

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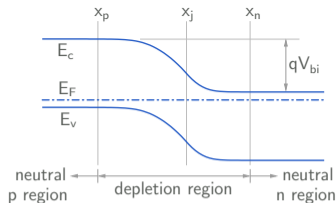
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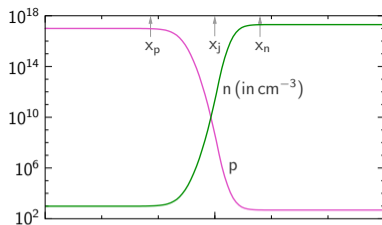
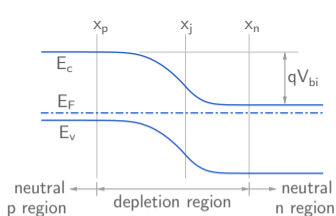


pn junction in equilibrium

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pn junction in equilibrium

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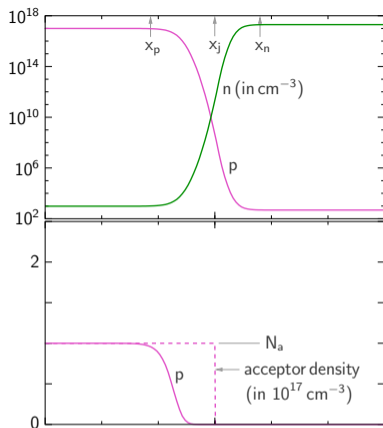
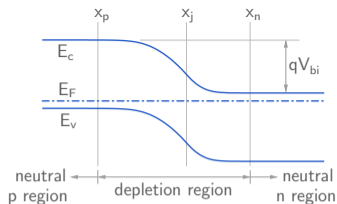
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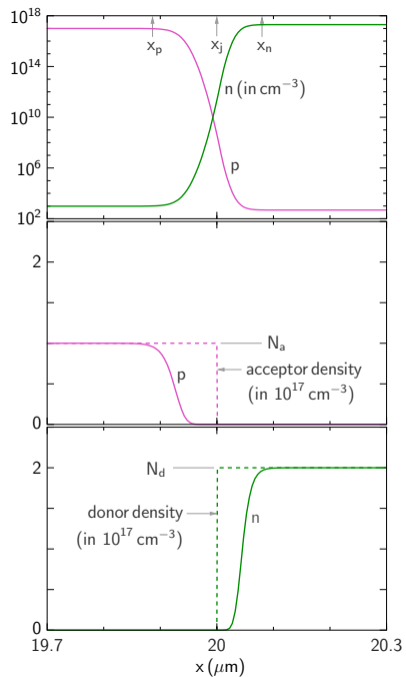
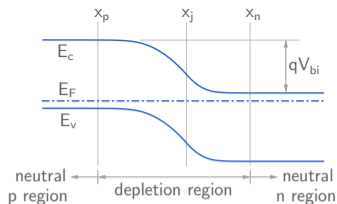
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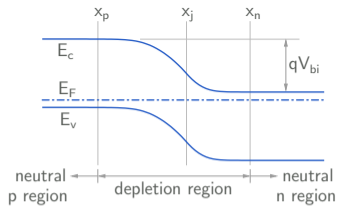
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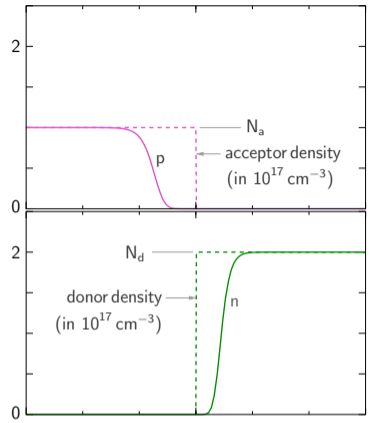
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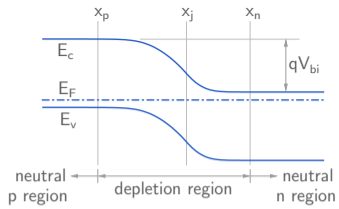
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$$\rho(x) = q [N_d^+(x) + p(x) - N_a^-(x) - n(x)]$$

$$\approx q [(N_d(x) - n(x)) - (N_a(x) - p(x))].$$



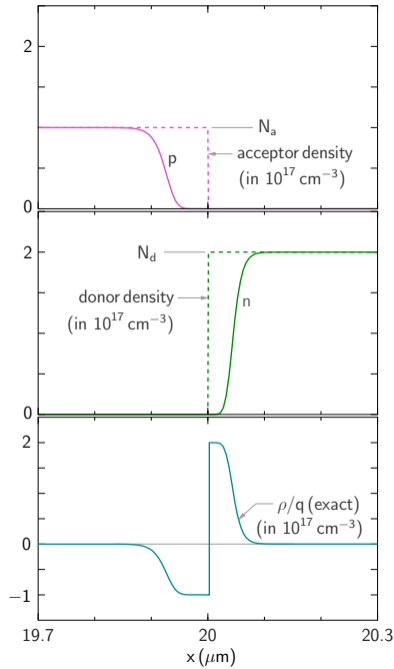
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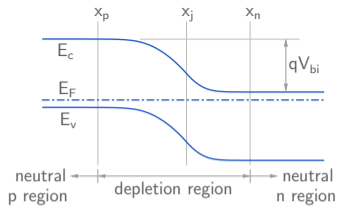
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pn junction in equilibrium

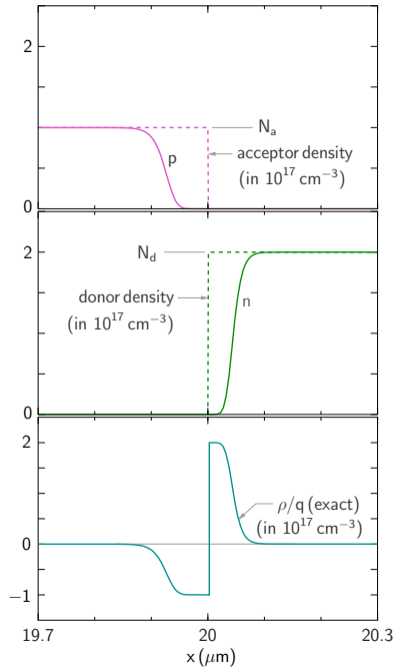


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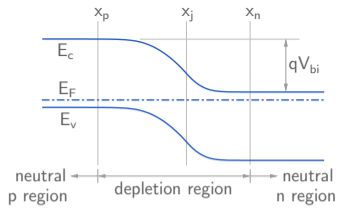
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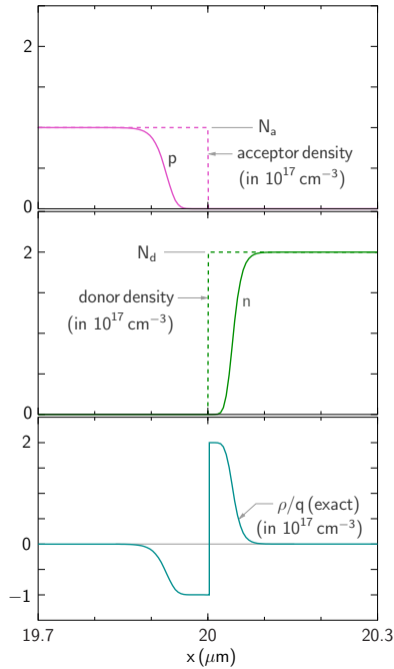
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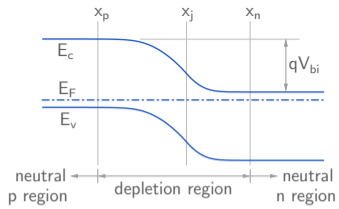
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* Within the depletion region, both n and p are small, i.e., this region is depleted of carriers \rightarrow "depletion region".



pn junction in equilibrium



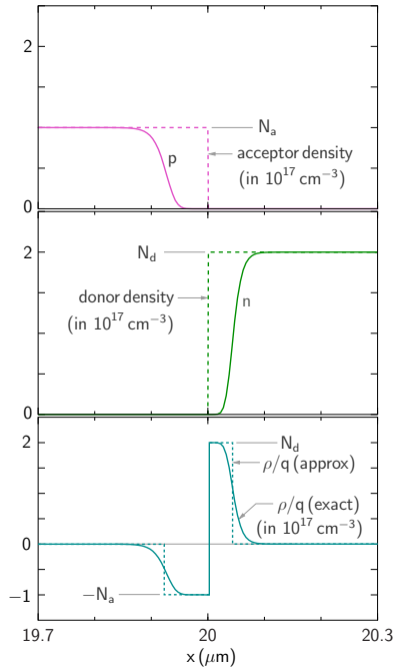
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$$\rho(x) = q [N_d^+(x) + p(x) - N_a^-(x) - n(x)]$$

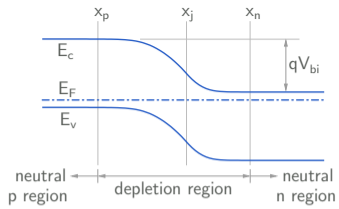
$$\approx q [(N_d(x) - n(x)) - (N_a(x) - p(x))].$$

* $\rho = 0$ in the neutral regions.

* Within the depletion region, both n and p are small, i.e., this region is depleted of carriers \rightarrow "depletion region".



pn junction in equilibrium



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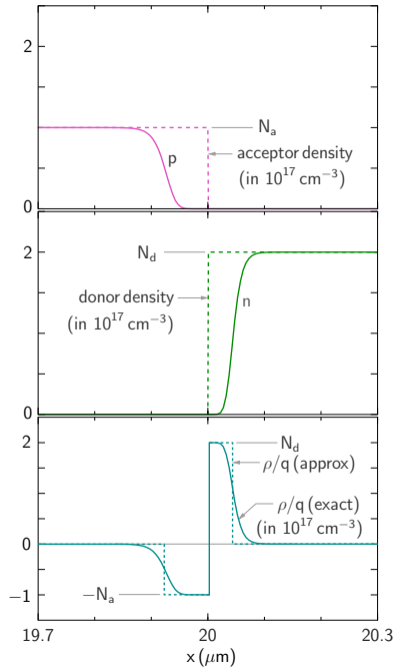
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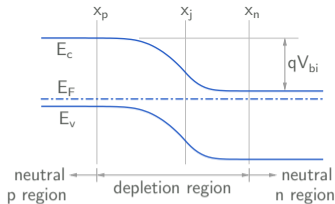
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pn junction in equilibrium



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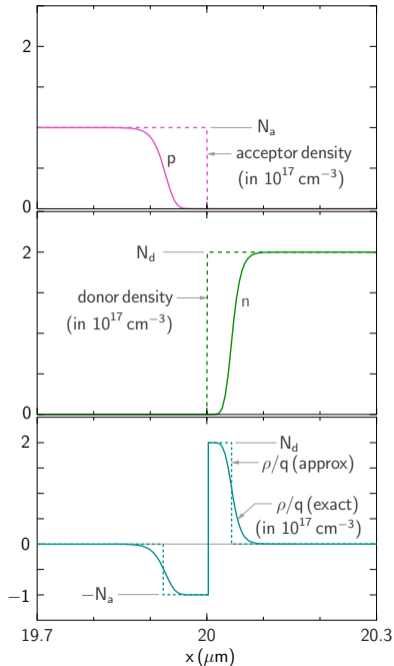
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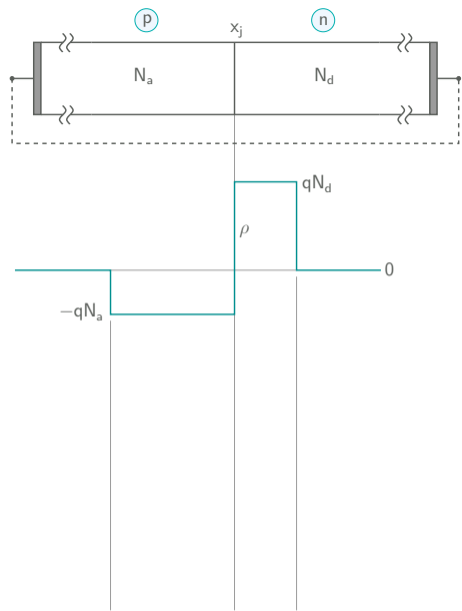
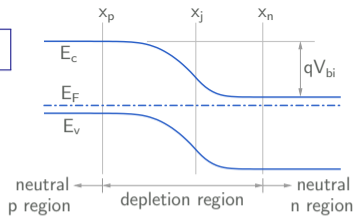
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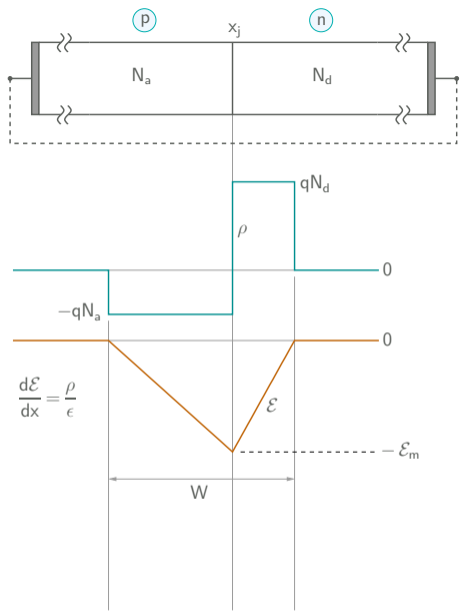
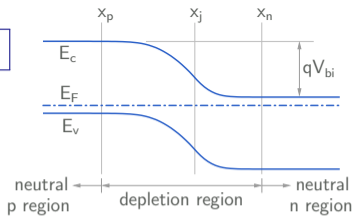
* Since the depletion region has non-zero charge density, it is also called "space charge region."



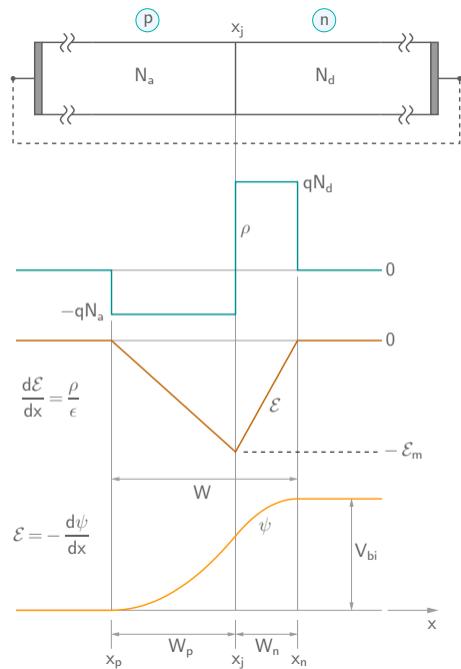
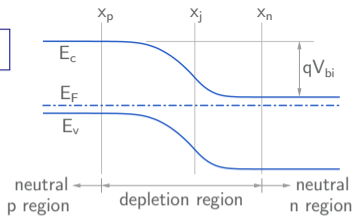
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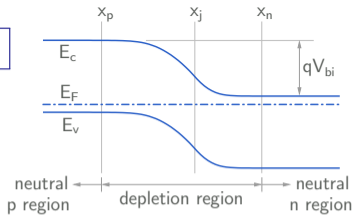
pn junction in equilibrium



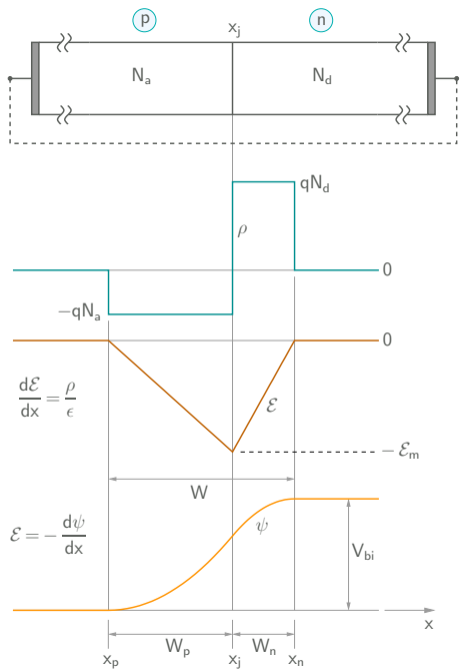
pn junction in equilibrium



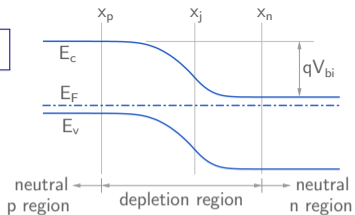
pn junction in equilibrium



* Built-in voltage V_{bi} :



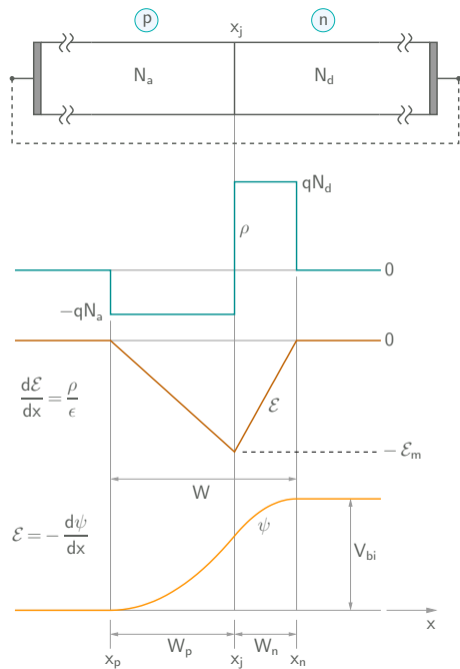
pn junction in equilibrium



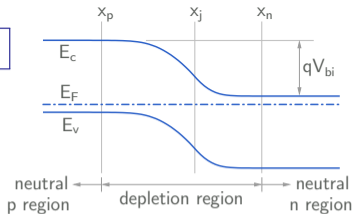
* Built-in voltage V_{bi} :

$$p(x) = N_v \exp \left[-\frac{E_F - E_v(x)}{kT} \right], \quad n(x) = N_c \exp \left[-\frac{E_c(x) - E_F}{kT} \right].$$

$$\rightarrow \frac{p(x_n)}{p(x_p)} = \exp \left[-\frac{E_v(x_p) - E_v(x_n)}{kT} \right].$$



pn junction in equilibrium

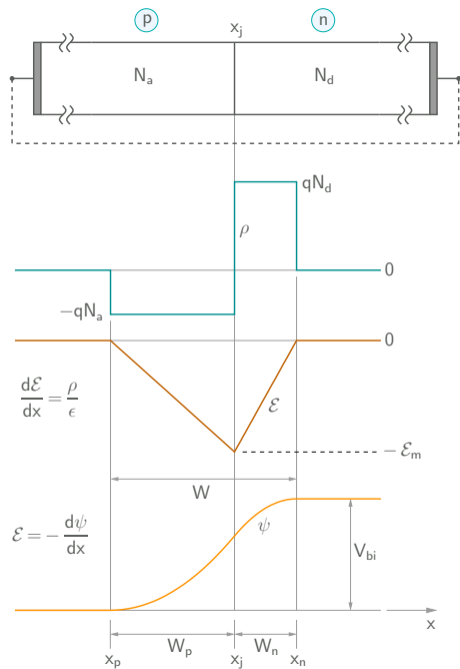


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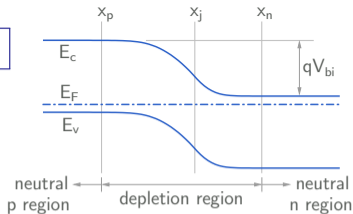
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$$\rightarrow \frac{p(x_n)}{p(x_p)} = \exp \left[-\frac{E_v(x_p) - E_v(x_n)}{kT} \right].$$

$$\rightarrow \frac{p(x_n)n(x_n)}{p(x_p)n(x_p)} = \frac{n_i^2}{N_a N_d} = \exp \left(-\frac{qV_{bi}}{kT} \right).$$



pn junction in equilibrium



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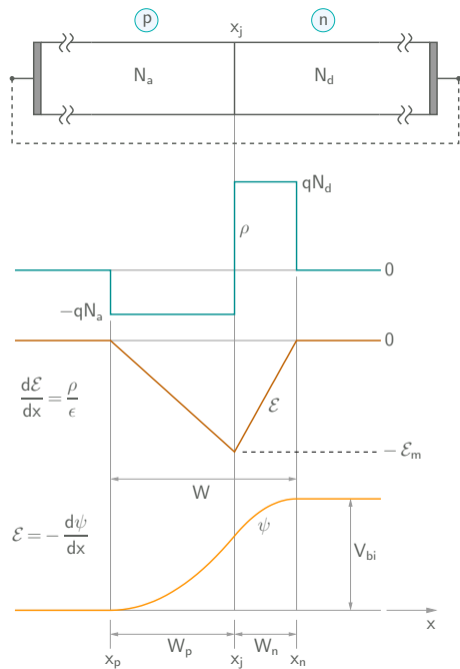
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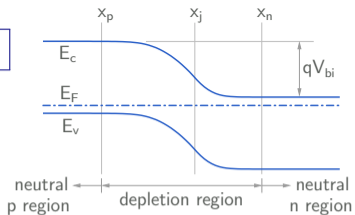
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The built-in voltage V_{bi} is therefore given by

$$V_{bi} = \frac{kT}{q} \log \left(\frac{N_a N_d}{n_i^2} \right)$$

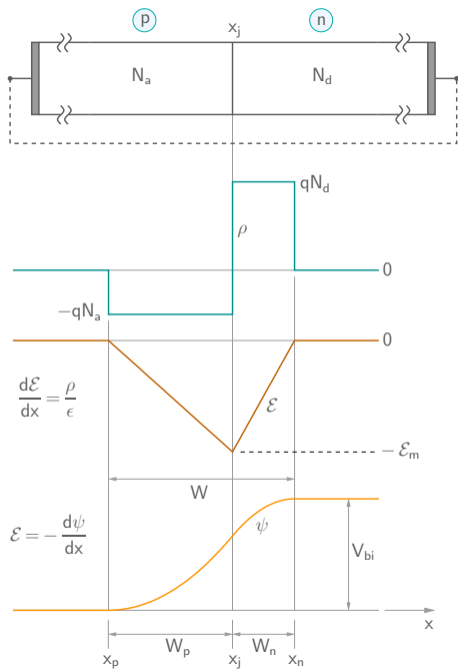


pn junction in equilibrium

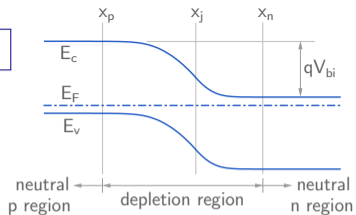


Example:

For a silicon pn junction with $N_a = 5 \times 10^{17} \text{ cm}^{-3}$, $N_d = 10^{17} \text{ cm}^{-3}$, compute V_{bi} at $T = 300 \text{ K}$. ($n_i = 1.5 \times 10^{10} \text{ cm}^{-3}$ at 300 K.)



pn junction in equilibrium

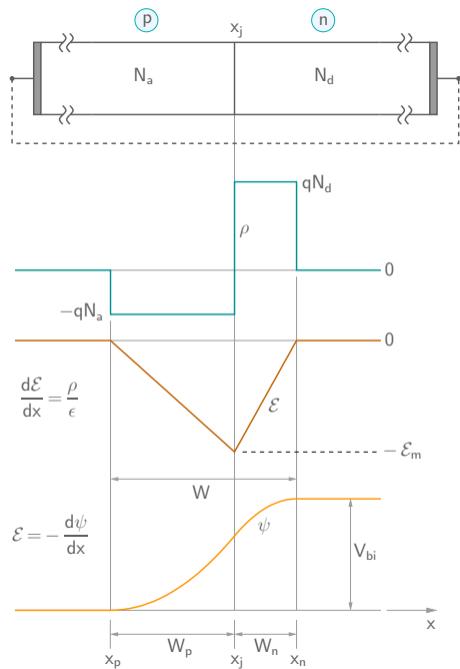


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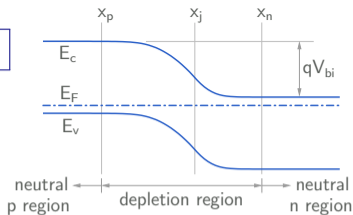
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Solution:

$$V_{bi} = \frac{kT}{q} \log \frac{N_a N_d}{n_i^2}$$



pn junction in equilibrium



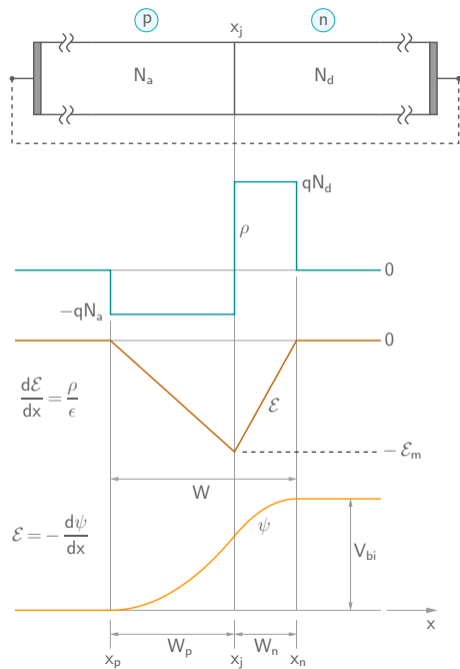
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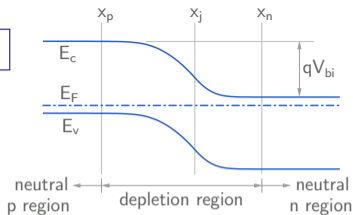
Solution:

$$V_{bi} = \frac{kT}{q} \log \frac{N_a N_d}{n_i^2}$$

$$= (0.0259 \text{ V}) \log \frac{(5 \times 10^{17})(10^{17})}{(1.5 \times 10^{10})^2}$$



pn junction in equilibrium

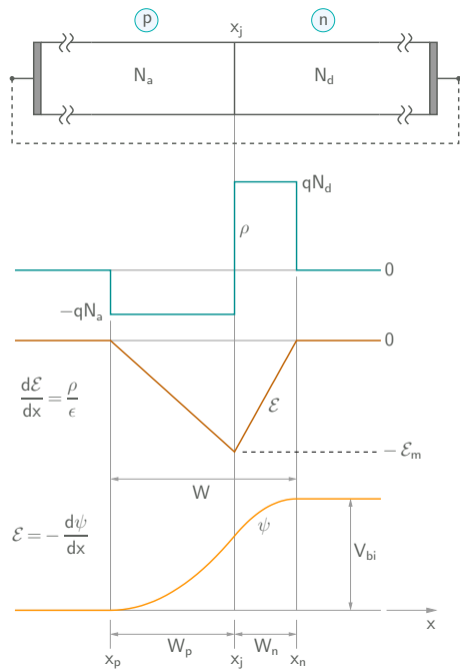


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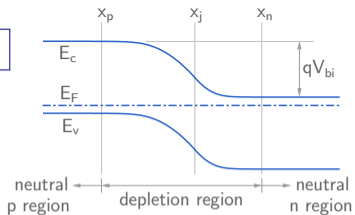
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Solution:

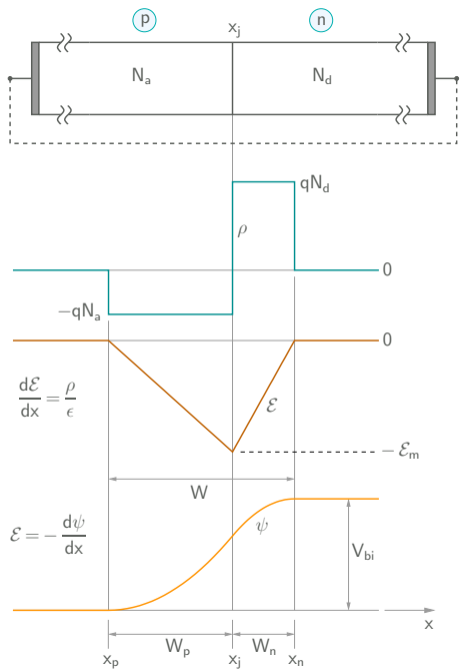
$$\begin{aligned}
 V_{bi} &= \frac{kT}{q} \log \frac{N_a N_d}{n_i^2} \\
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 &= 0.86 \text{ V}
 \end{aligned}$$



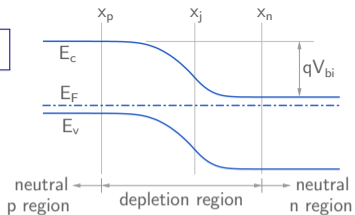
pn junction in equilibrium



Electric field $\mathcal{E}(x)$:

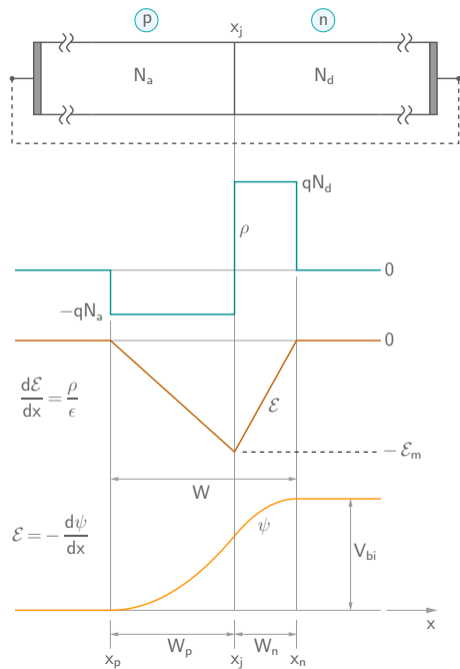


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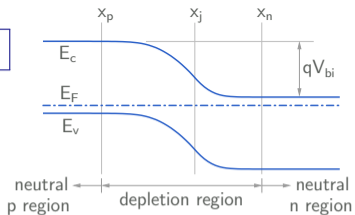


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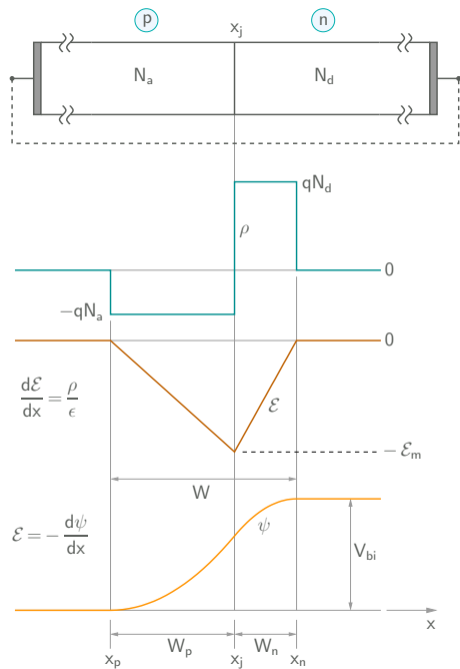
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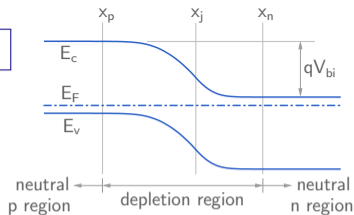
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* In the depletion region, $\int_{x_p}^{x_n} d\mathcal{E} = \int_{x_p}^{x_n} \frac{\rho}{\epsilon} dx$.



pn junction in equilibrium



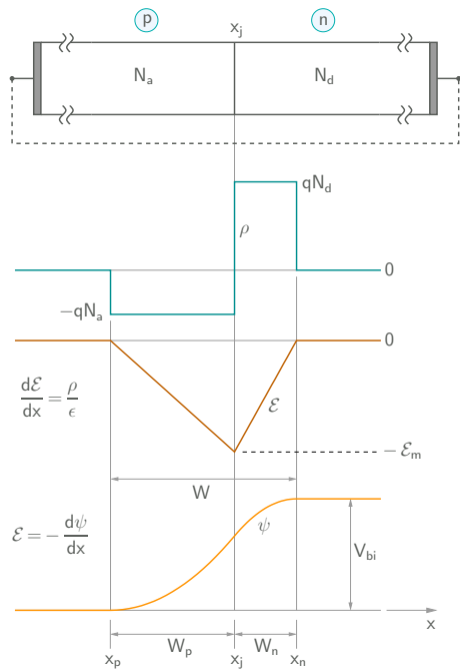
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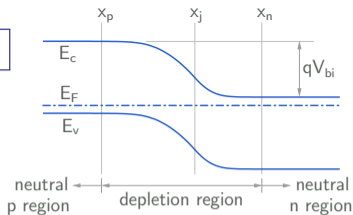
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which means the area under the ρ versus x curve must be zero.



pn junction in equilibrium



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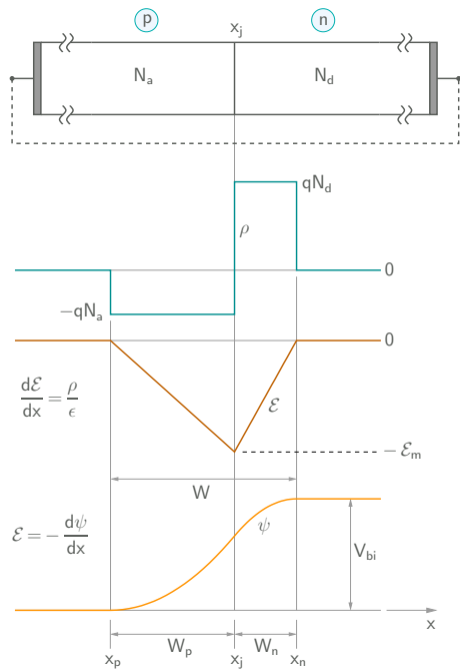
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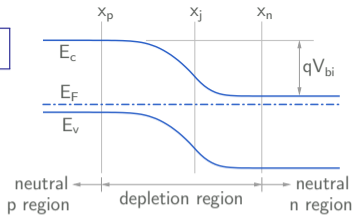
which means the area under the ρ versus x curve must be zero.

i.e., $N_a W_p = N_d W_n \rightarrow \frac{W_p}{W_n} = \frac{N_d}{N_a}$.

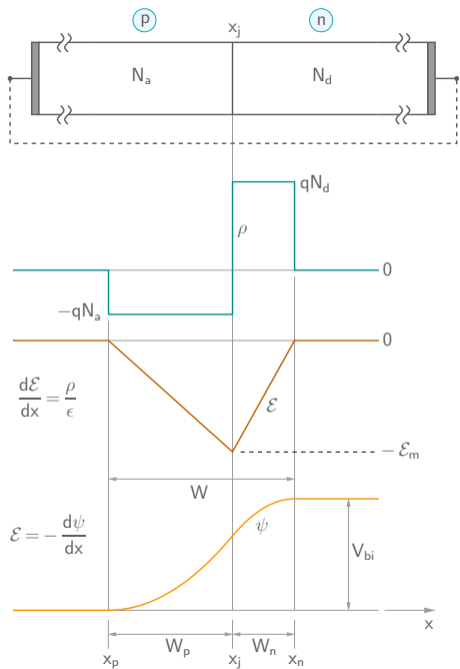
→ The depletion width is larger on the lightly doped side.



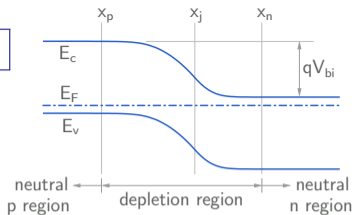
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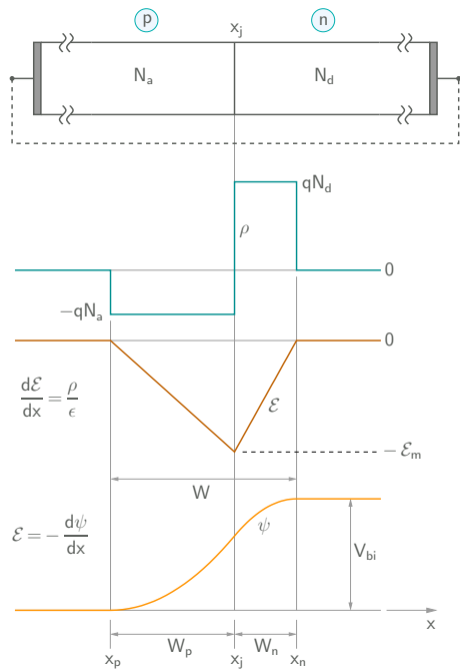


pn junction in equilibrium

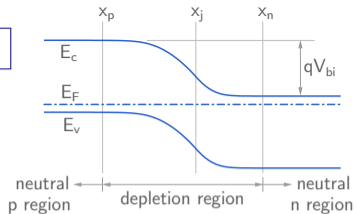


Electric field $\mathcal{E}(x)$:

* Since ρ is piecewise constant, \mathcal{E} must be piecewise linear.



pn junction in equilibrium

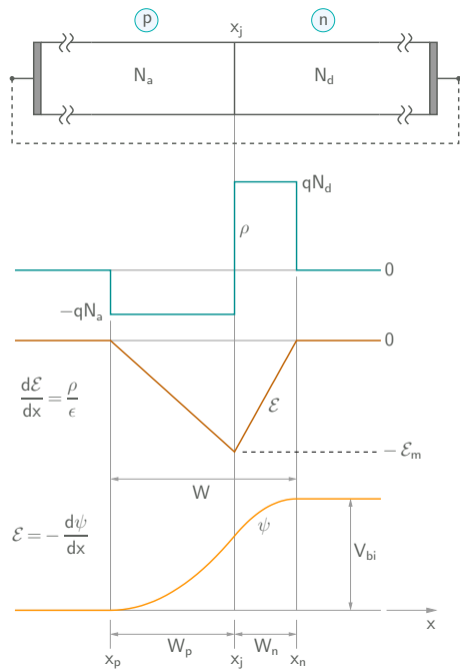


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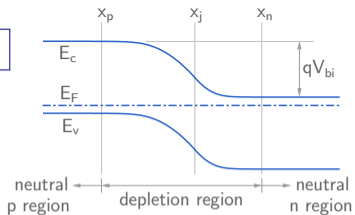
- * Since ρ is piecewise constant, \mathcal{E} must be piecewise linear.
- * The maximum value (magnitude) of \mathcal{E} occurs at $x = x_j$.

$$\int_{x_p}^{x_j} d\mathcal{E} = \frac{1}{\epsilon} \int_{x_p}^{x_j} \rho dx \rightarrow -\mathcal{E}_m - 0 = \frac{1}{\epsilon} (-qN_a W_p)$$

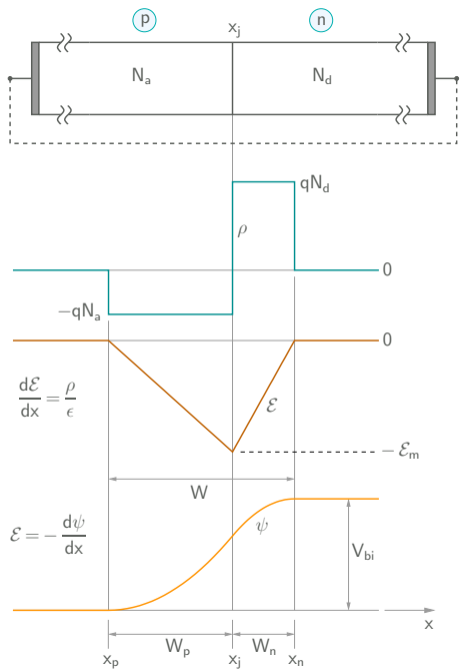
$$\rightarrow \mathcal{E}_m = \frac{qN_a W_p}{\epsilon} = \frac{qN_d W_n}{\epsilon} \quad \therefore N_a W_p = N_d W_n.$$



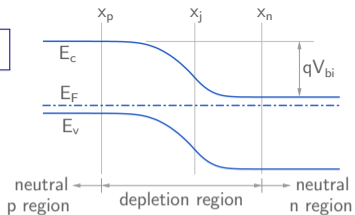
pn junction in equilibrium



Potential $\psi(x)$:

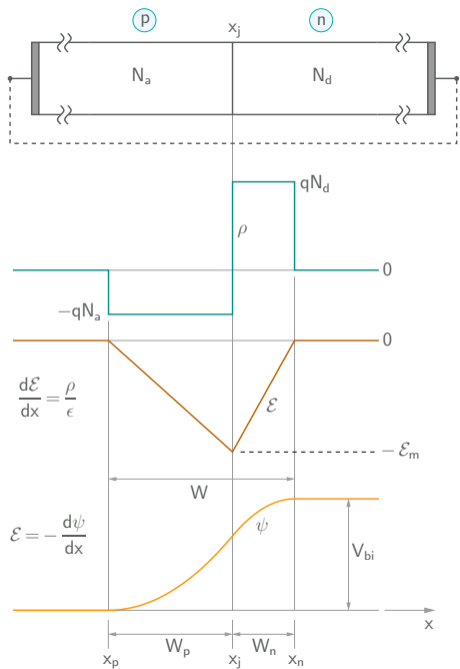


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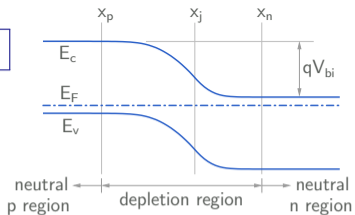


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$$* \quad x_p < x < x_j: \quad \frac{d\mathcal{E}}{dx} = -\frac{qN_a^-}{\epsilon} \approx -\frac{qN_a}{\epsilon} \rightarrow \mathcal{E}(x) = -\frac{qN_a}{\epsilon}x + k_1.$$



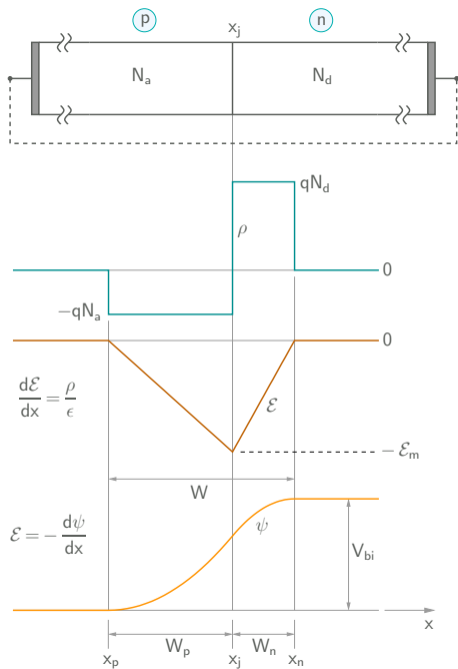
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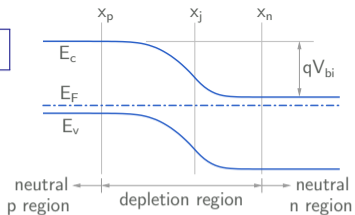
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$$\text{Since } \mathcal{E} = 0 \text{ at } x = x_p, \text{ we get } \mathcal{E}(x) = -\frac{qN_a}{\epsilon}(x - x_p).$$



pn junction in equilibrium

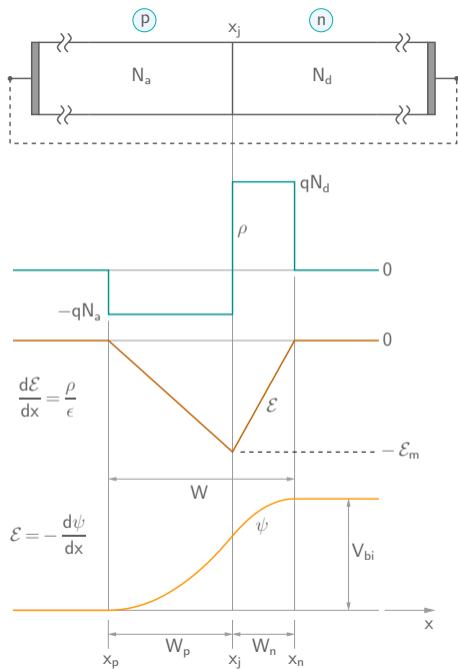


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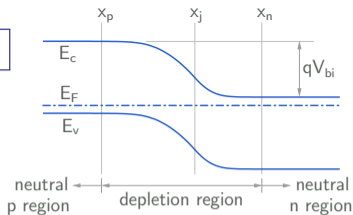
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$$\rightarrow \psi(x) = -\int \mathcal{E} dx = \frac{qN_a}{\epsilon} \left[\frac{x^2}{2} - x_p x \right] + k_2.$$



pn junction in equilibrium



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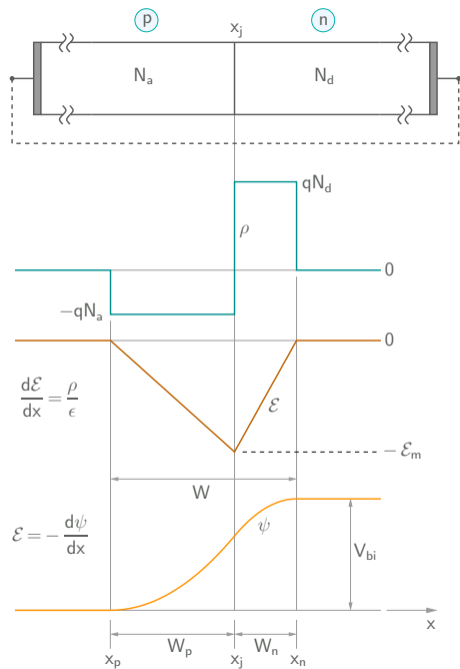
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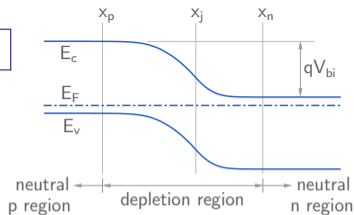
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Taking $\psi(x_p) = 0$, we can find k_2 .

$$\rightarrow \psi(x) = \frac{qN_a}{2\epsilon}(x - x_p)^2.$$



pn junction in equilibrium



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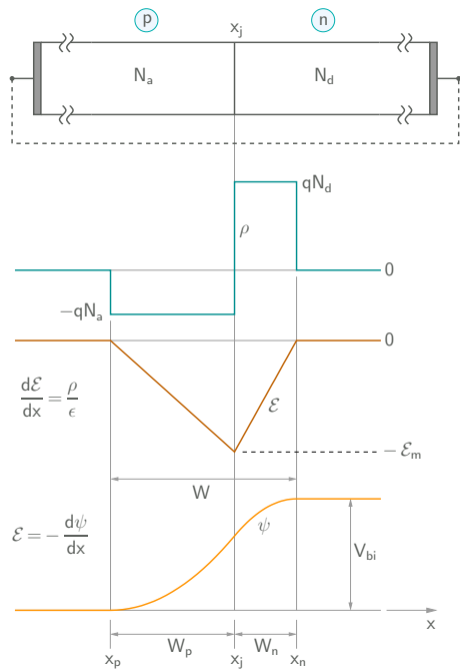
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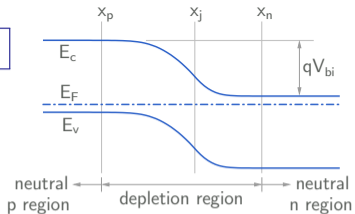
$$\rightarrow \psi(x) = \frac{qN_a}{2\epsilon}(x - x_p)^2.$$

If x_j is taken as 0, i.e., $x \leftarrow (x - x_j)$, we get

$$\psi(x) = \frac{qN_a}{2\epsilon}(x + W_p)^2.$$

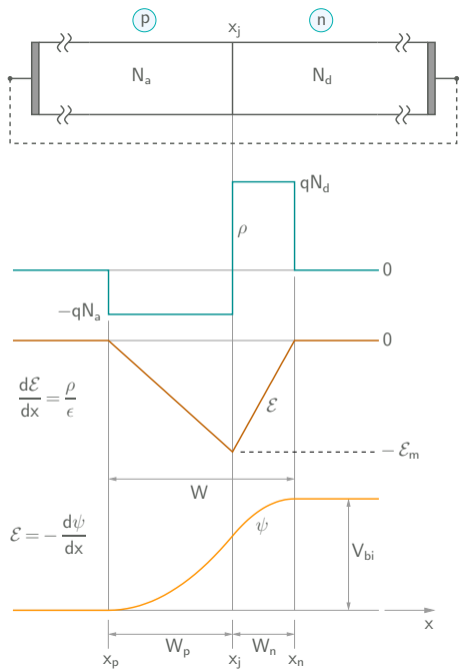


pn junction in equilibrium

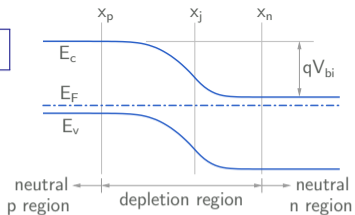


Potential $\psi(x)$:

* $x_j < x < x_n$:



pn junction in equilibrium

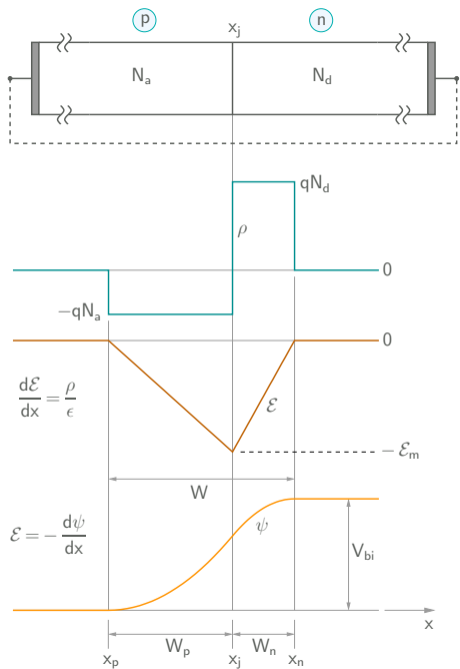


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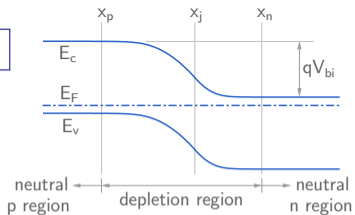
* $x_j < x < x_n$:

For convenience, let us take $x_j = 0 \rightarrow x_p = -W_p, x_n = W_n$.

$$\frac{d\mathcal{E}}{dx} = \frac{qN_d^+}{\epsilon} \approx \frac{qN_d}{\epsilon} \rightarrow \mathcal{E}(x) = \frac{qN_d}{\epsilon}x + k_3.$$



pn junction in equilibrium



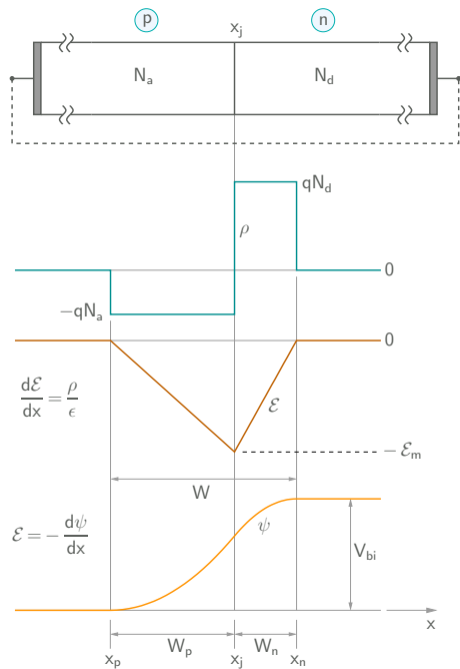
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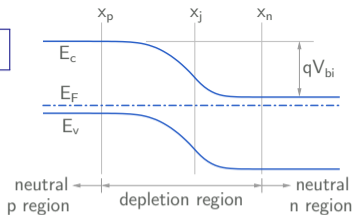
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$$\text{Since } \mathcal{E} = 0 \text{ at } x = W_n, \text{ we get } \mathcal{E}(x) = \frac{qN_d}{\epsilon}(x - W_n).$$



pn junction in equilibrium



Potential $\psi(x)$:

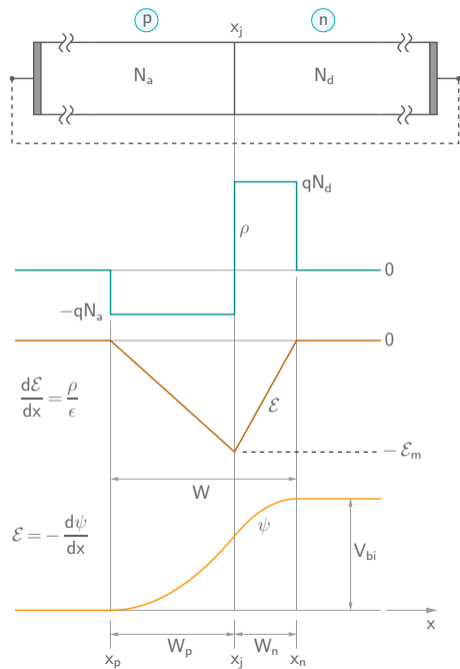
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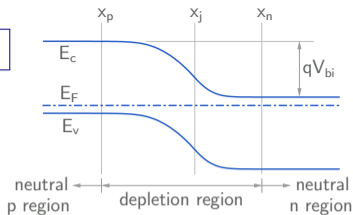
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Since $\mathcal{E} = 0$ at $x = W_n$, we get $\mathcal{E}(x) = \frac{qN_d}{\epsilon}(x - W_n)$.

$$\rightarrow \psi(x) = - \int \mathcal{E} dx = - \frac{qN_d}{\epsilon} \left[\frac{x^2}{2} - W_n x \right] + k_4.$$



pn junction in equilibrium



Potential $\psi(x)$:

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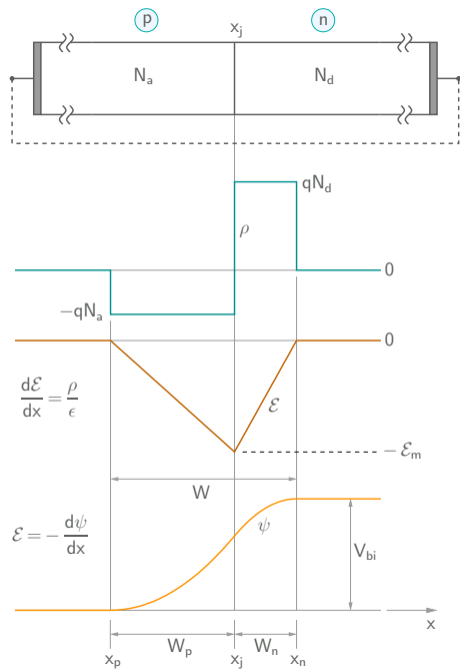
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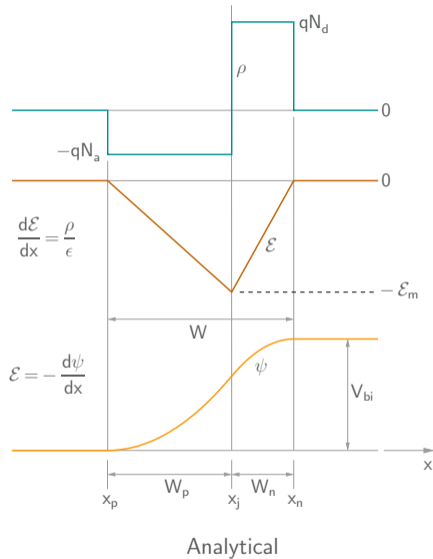
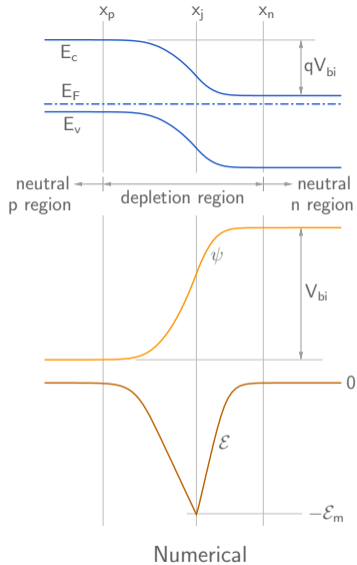
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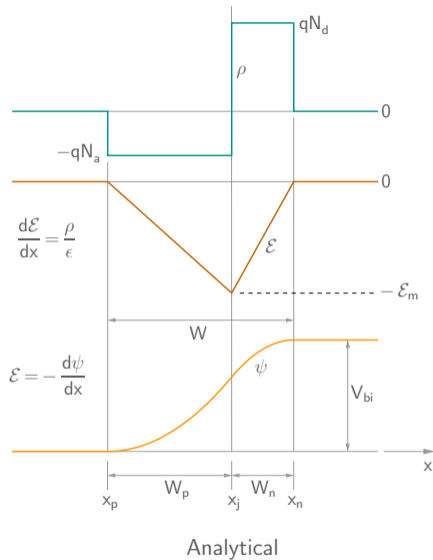
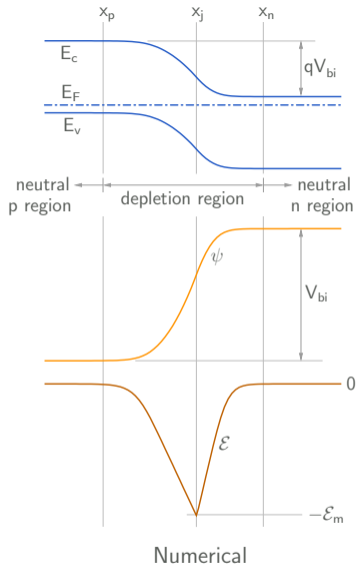
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We can find k_4 using continuity of ψ at $x = 0$.

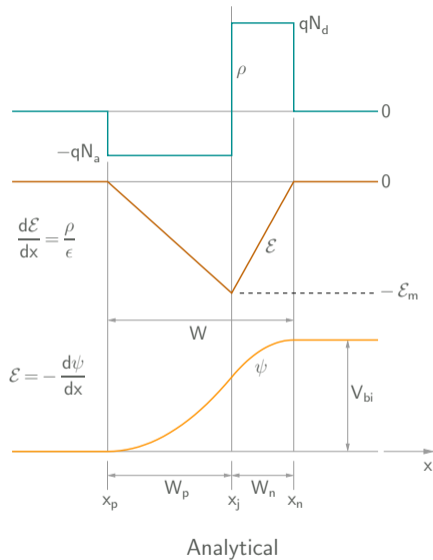
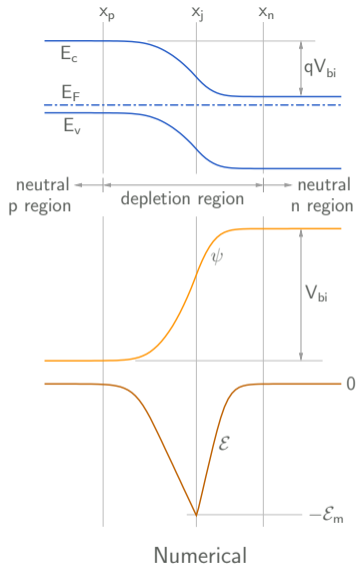
$$\rightarrow \psi(x) = \frac{qN_d}{\epsilon} \left[W_n x - \frac{x^2}{2} \right] + \frac{qN_a}{2\epsilon} W_p^2.$$





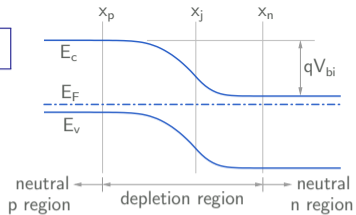


* *pn* junction in equilibrium: The band diagram is consistent with Poisson's equation.

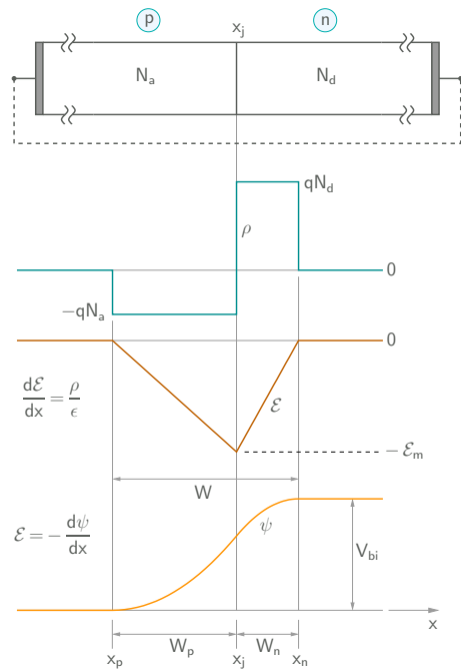


* *pn* junction in equilibrium: Depletion approximation agrees well with numerical results.

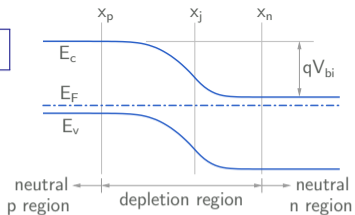
pn junction in equilibrium



Depletion region width W :

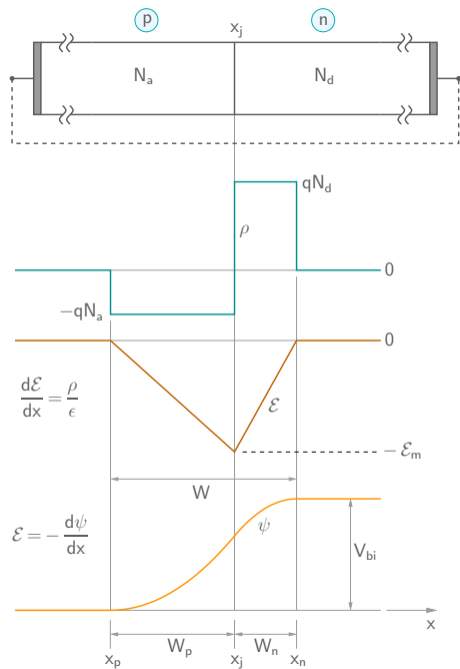


pn junction in equilibrium

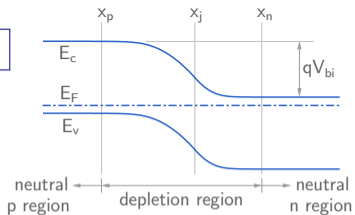


Depletion region width W :

The built-in voltage V_{bi} is given by the area under the $\mathcal{E}(x)$ curve.



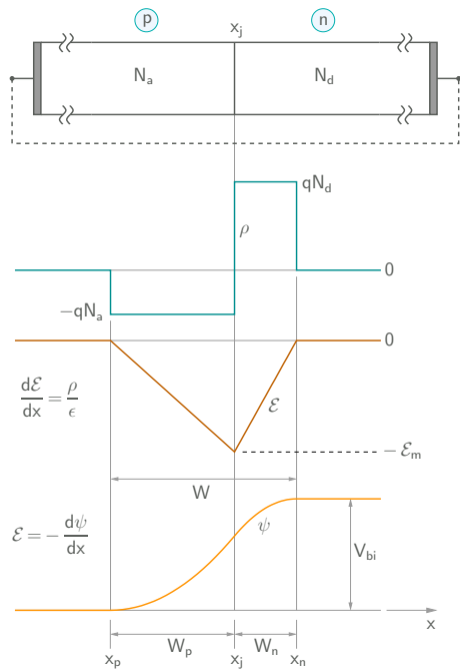
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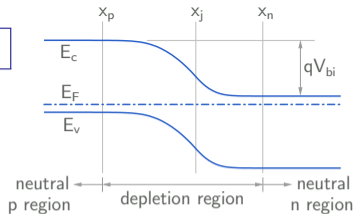
Depletion region width W :

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$$V_{bi} = \frac{1}{2} \mathcal{E}_m W_p + \frac{1}{2} \mathcal{E}_m W_n = \frac{1}{2} \mathcal{E}_m W = \frac{1}{2} \frac{qN_a W_p}{\epsilon} W.$$



pn junction in equilibrium



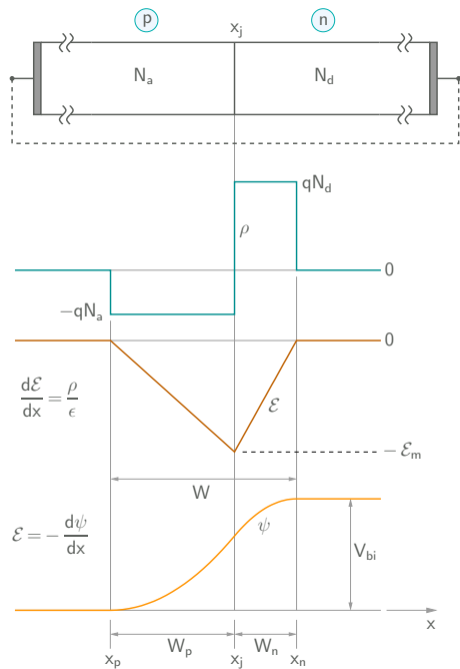
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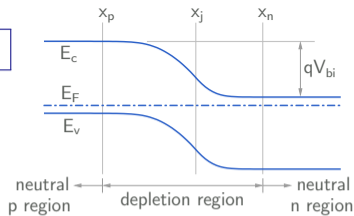
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Since $W_n + W_p = W$ and $W_n N_d = W_p N_a$, we get

$$W_n = \frac{N_a}{N_a + N_d} W, \quad W_p = \frac{N_d}{N_a + N_d} W.$$



pn junction in equilibrium



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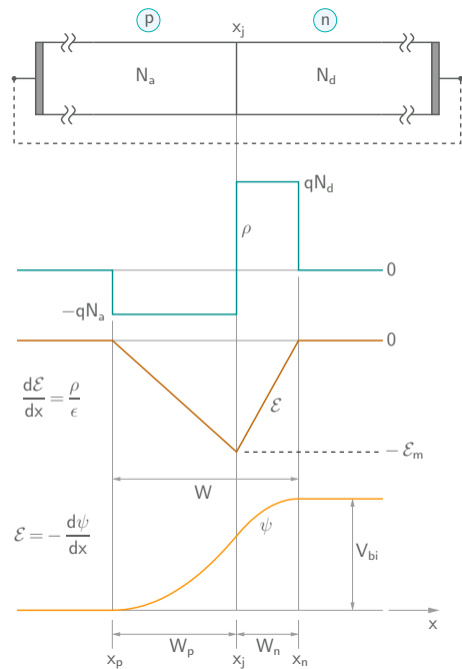
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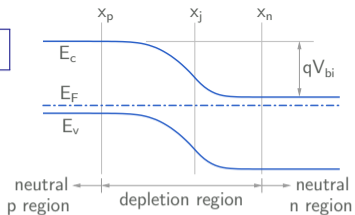
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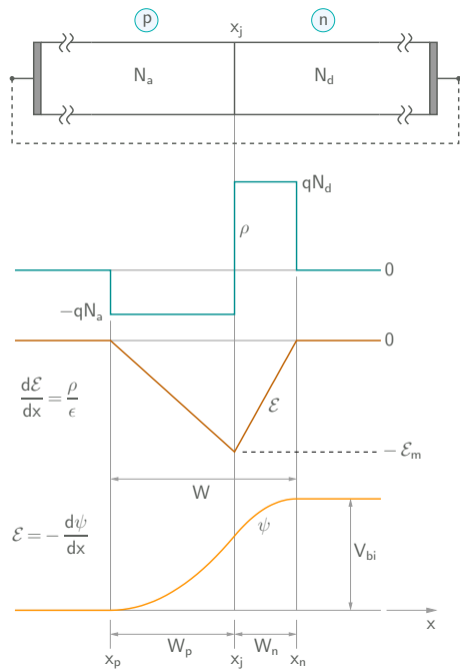
$$\rightarrow V_{bi} = \frac{1}{2} \frac{q}{\epsilon} \frac{N_a N_d}{N_a + N_d} W^2, \quad \text{i.e., } W = \sqrt{\frac{2\epsilon}{q} \left(\frac{N_a + N_d}{N_a N_d} \right) V_{bi}}.$$



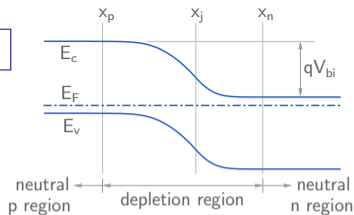
pn junction in equilibrium



For an abrupt, uniformly doped silicon *pn* junction, $N_a = 5 \times 10^{17} \text{ cm}^{-3}$. Compute V_{bi} , W , W_n , W_p , and \mathcal{E}_m for $N_d = 10^{16}, 10^{17}, 5 \times 10^{17}, 10^{18}$, and $5 \times 10^{18} \text{ cm}^{-3}$ ($T = 300 \text{ K}$).



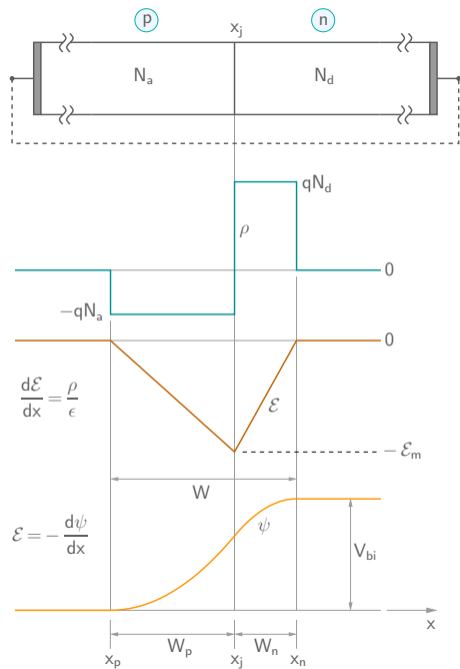
pn junction in equilibrium



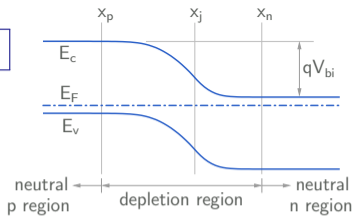
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Solution:

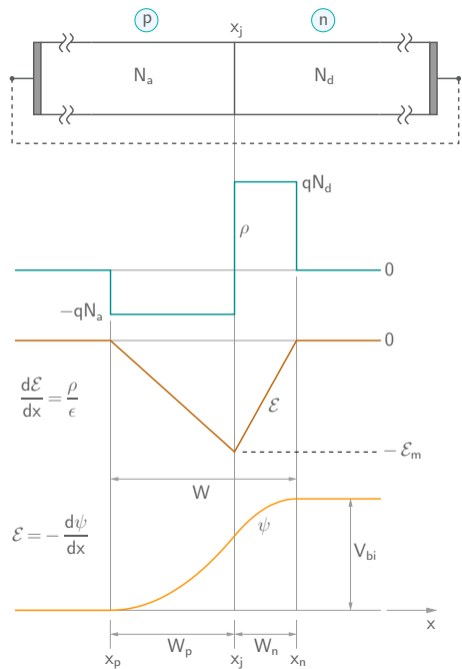
$$\begin{aligned}
 V_{bi} &= V_T \log \frac{N_a N_d}{n_i^2} \\
 &= 0.0259 \text{ V} \times \log \frac{(5 \times 10^{17})(1 \times 10^{16})}{(1.5 \times 10^{10})^2} \\
 &= 0.8 \text{ V}.
 \end{aligned}$$



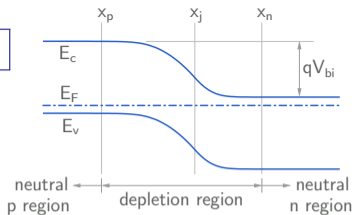
pn junction in equilibrium



$$W = \sqrt{\frac{2\epsilon_r\epsilon_0}{q} \left(\frac{N_a + N_d}{N_a N_d} \right) V_{bi}}$$

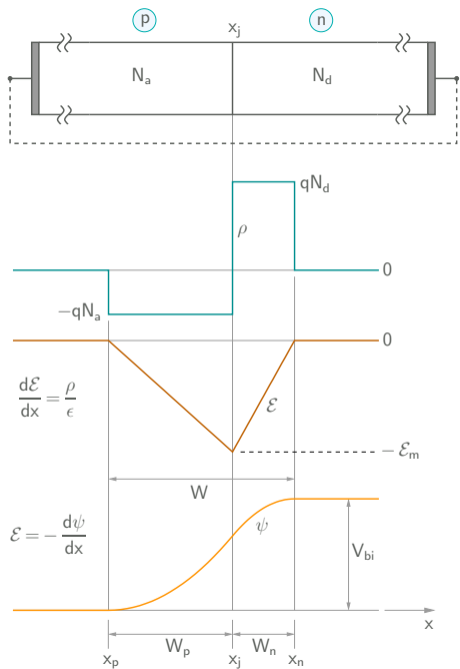


pn junction in equilibrium

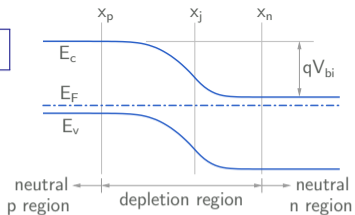


$$W = \sqrt{\frac{2\epsilon_r\epsilon_0}{q} \left(\frac{N_a + N_d}{N_a N_d} \right) V_{bi}}$$

$$= \sqrt{\frac{2 \times 11.7 \times 8.85 \times 10^{-14}}{1.6 \times 10^{-19}} \left(\frac{5 \times 10^{17} + 1 \times 10^{16}}{(5 \times 10^{17})(1 \times 10^{16})} \right) (0.8)}$$



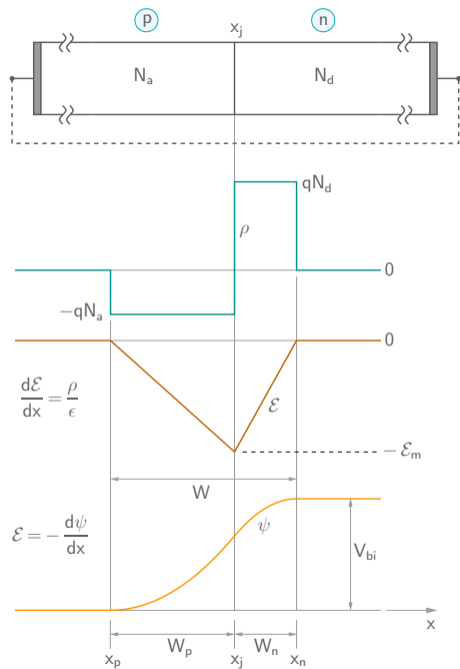
pn junction in equilibrium



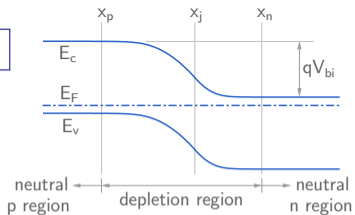
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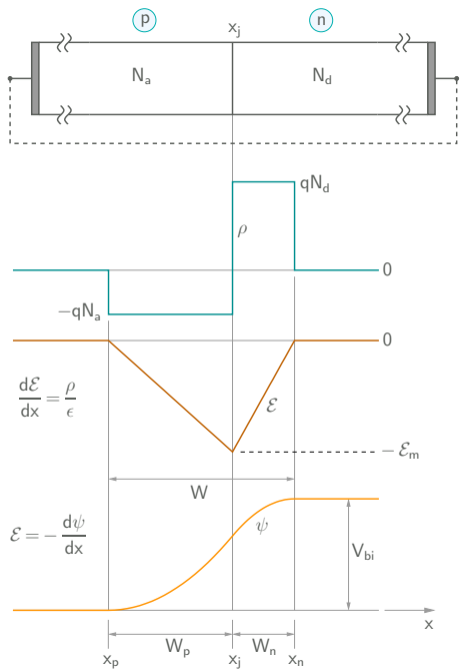
$$= 3.24 \times 10^{-5} \text{ cm}$$



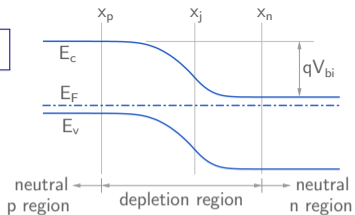
pn junction in equilibrium



$$\begin{aligned}
 W &= \sqrt{\frac{2\epsilon_r\epsilon_0}{q} \left(\frac{N_a + N_d}{N_a N_d} \right) V_{bi}} \\
 &= \sqrt{\frac{2 \times 11.7 \times 8.85 \times 10^{-14}}{1.6 \times 10^{-19}} \left(\frac{5 \times 10^{17} + 1 \times 10^{16}}{(5 \times 10^{17})(1 \times 10^{16})} \right) (0.8)} \\
 &= 3.24 \times 10^{-5} \text{ cm} \\
 &= 0.324 \mu\text{m}.
 \end{aligned}$$



pn junction in equilibrium



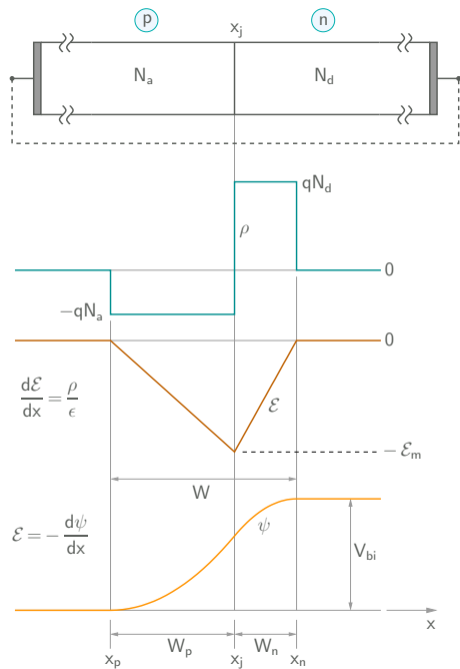
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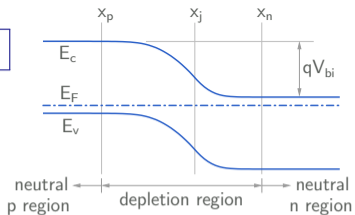
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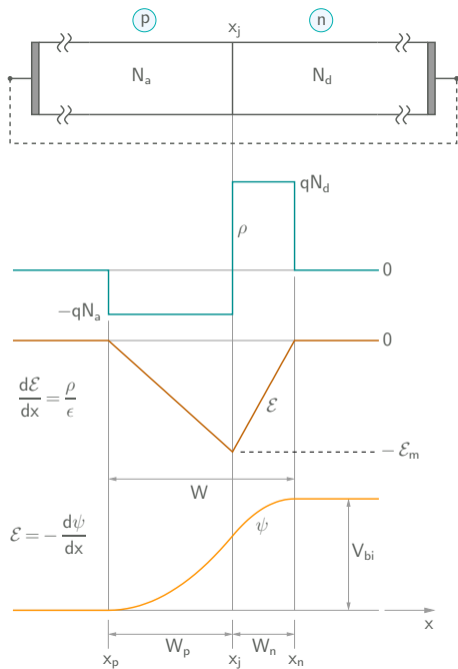
$$\text{Units: } \sqrt{\frac{\text{F/cm}}{\text{Coul}} \times \frac{\text{cm}^{-3}}{\text{cm}^{-6}} \times \text{V}} = \text{cm}$$



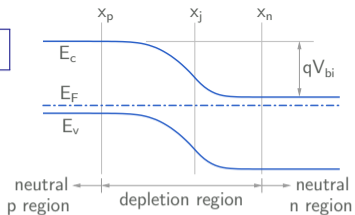
pn junction in equilibrium



$$W_n = \frac{N_a}{N_a + N_d} W = 0.318 \mu\text{m}, \quad W_p = \frac{N_d}{N_a + N_d} W = 0.006 \mu\text{m}.$$

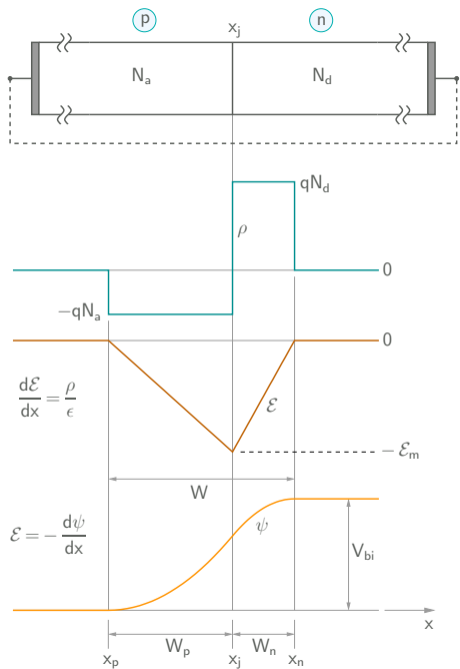


pn junction in equilibrium

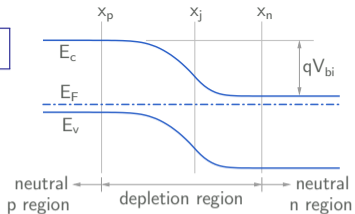


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$$\mathcal{E}_m = \frac{qN_d}{\epsilon} W_n \quad \text{or} \quad \frac{qN_a}{\epsilon} W_p$$



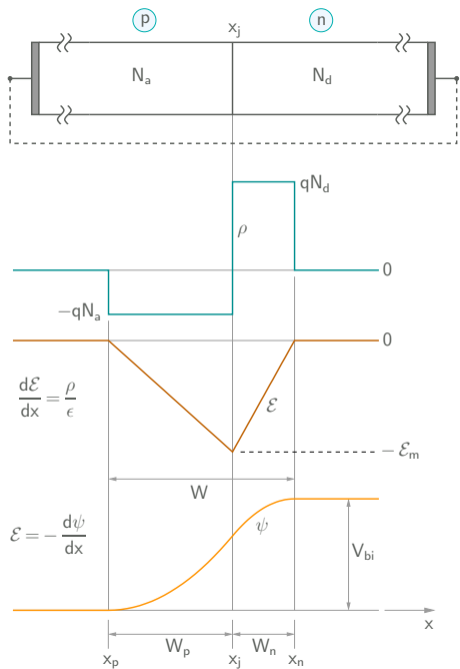
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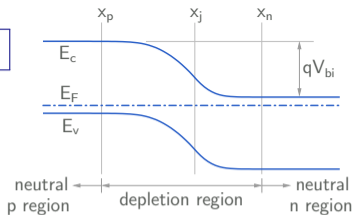
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$$\mathcal{E}_m = \frac{qN_d}{\epsilon} W_n \quad \text{or} \quad \frac{qN_a}{\epsilon} W_p$$

$$= \frac{1.6 \times 10^{-19} \text{ Coul} \times 10^{16} \text{ cm}^{-3}}{11.7 \times 8.85 \times 10^{-14} \text{ F/cm}} \times (3.18 \times 10^{-5} \text{ cm})$$

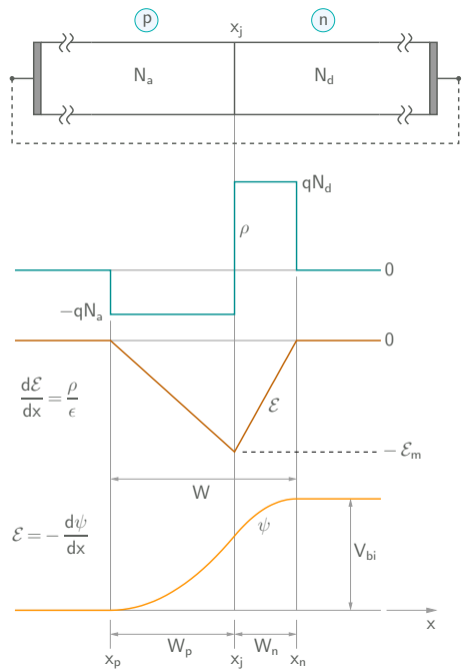


pn junction in equilibrium

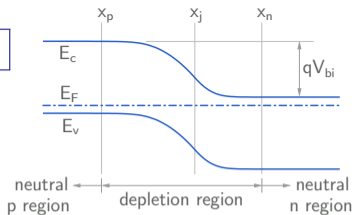


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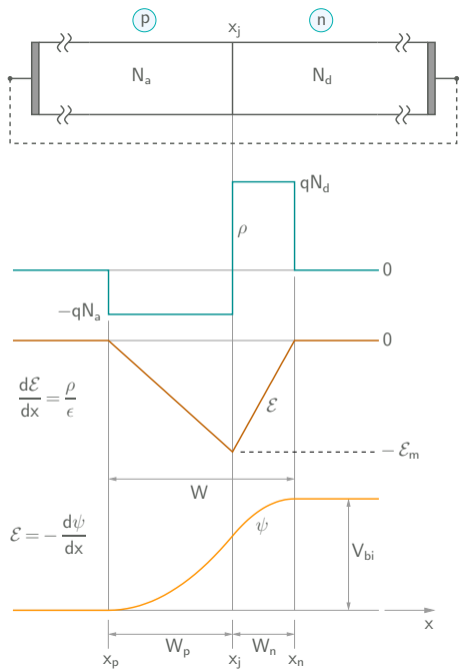


pn junction in equilibrium



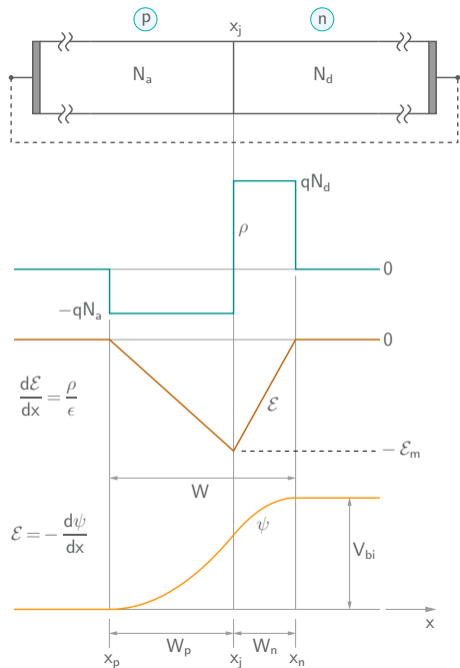
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Effect of N_d , with $N_a = 5 \times 10^{17} \text{ cm}^{-3}$ held fixed.
 (V_{bi} in Volts, W , W_n , W_p in μm , \mathcal{E}_m in kV/cm .)

$N_d (\text{cm}^{-3})$	V_{bi}	W	W_n	W_p	\mathcal{E}_m
1.0×10^{16}	0.80	0.324	0.318	0.006	49
1.0×10^{17}	0.86	0.115	0.096	0.019	148
5.0×10^{17}	0.90	0.068	0.034	0.034	263
1.0×10^{18}	0.92	0.060	0.020	0.040	307
5.0×10^{18}	0.96	0.052	0.004	0.047	366

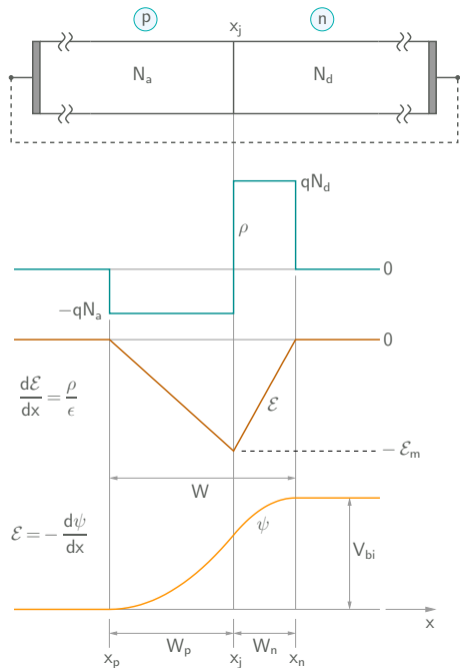


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 (V_{bi} in Volts, W , W_n , W_p in μm , \mathcal{E}_m in kV/cm .)

$N_d \text{ (cm}^{-3}\text{)}$	V_{bi}	W	W_n	W_p	\mathcal{E}_m
1.0×10^{16}	0.80	0.324	0.318	0.006	49
1.0×10^{17}	0.86	0.115	0.096	0.019	148
5.0×10^{17}	0.90	0.068	0.034	0.034	263
1.0×10^{18}	0.92	0.060	0.020	0.040	307
5.0×10^{18}	0.96	0.052	0.004	0.047	366

$$* V_{bi} = V_T \log \frac{N_a N_d}{n_i^2}, \quad W = \sqrt{\frac{2\epsilon}{q} \left(\frac{N_a + N_d}{N_a N_d} \right) V_{bi}},$$

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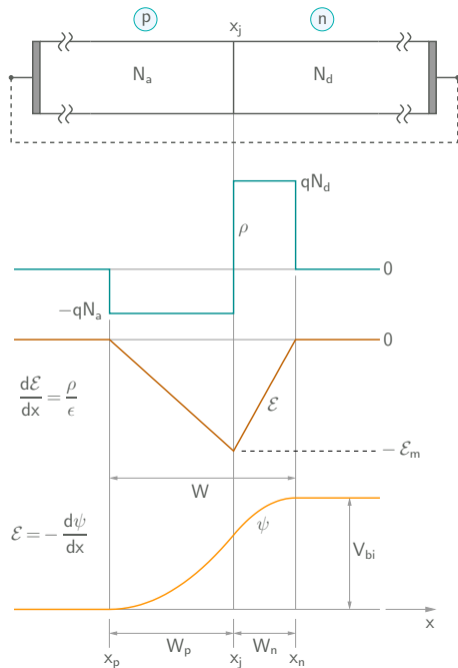
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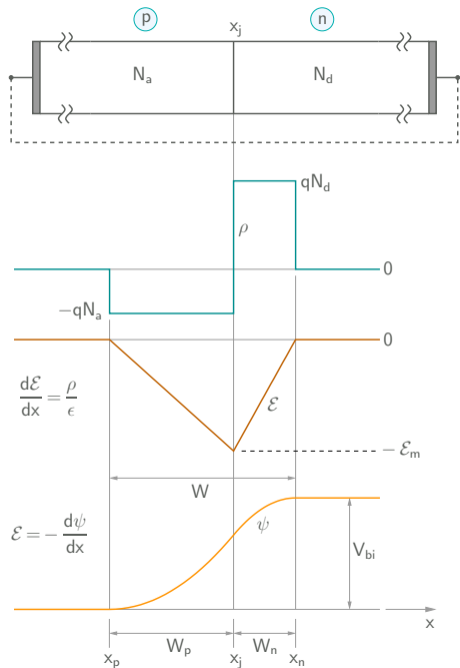
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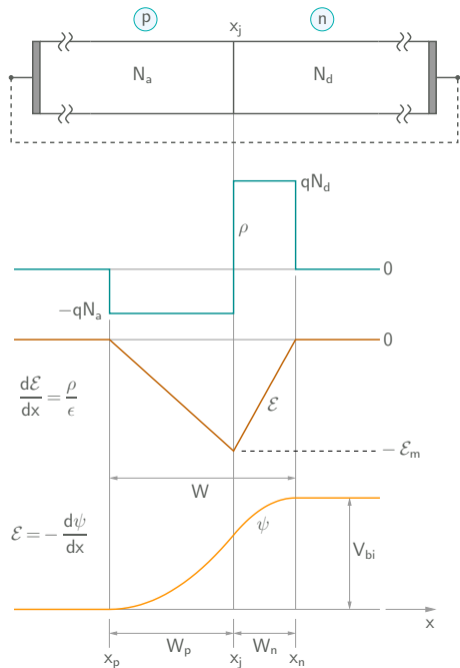
$N_a \ll N_d$ (n^+p junction):

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Effect of N_d , with $N_a = 5 \times 10^{17} \text{ cm}^{-3}$ held fixed.
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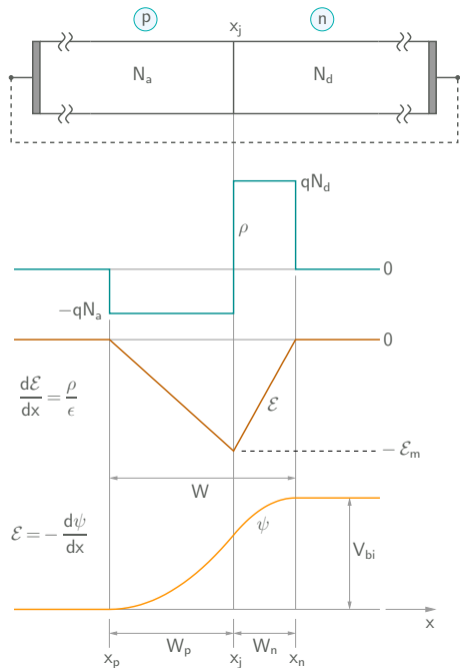
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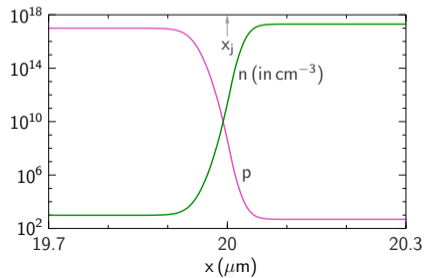
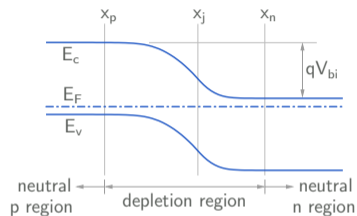
* For high doping densities such as 10^{18} cm^{-3} , degenerate statistics should be used for higher accuracy, i.e.,

$$n = N_c \frac{2}{\sqrt{\pi}} \mathcal{F}_{1/2}(\eta_c), \text{ with } \eta_c = \frac{E_F - E_c}{kT}, \text{ and}$$

$$p = N_v \frac{2}{\sqrt{\pi}} \mathcal{F}_{1/2}(\eta_v), \text{ with } \eta_v = \frac{E_v - E_F}{kT}.$$

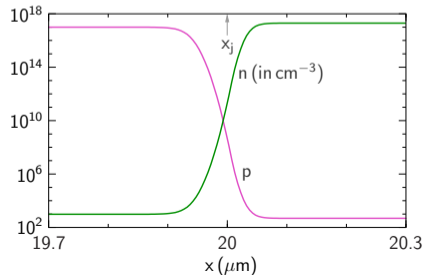
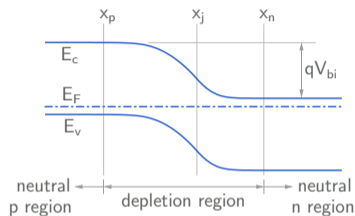
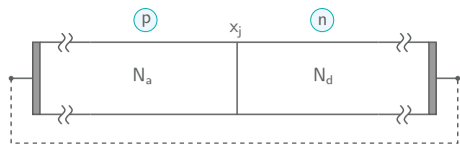


pn junction in equilibrium: current densities



pn junction in equilibrium: current densities

- * The diffusion currents can be expected to be substantial since there is a large change in n or p between the p -side and the n -side.

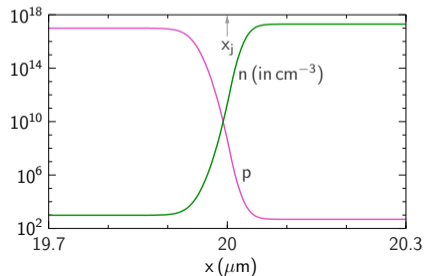
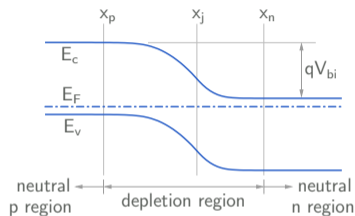


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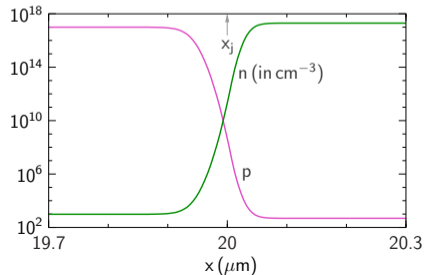
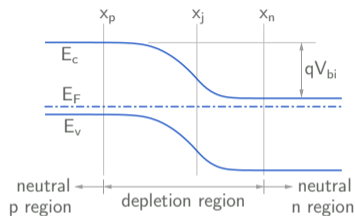
* In equilibrium, the drift and diffusion currents are equal and opposite for electrons as well as holes, i.e.,

$$J_n^{\text{diff}} = -J_n^{\text{drift}}, \quad J_p^{\text{diff}} = -J_p^{\text{drift}}.$$



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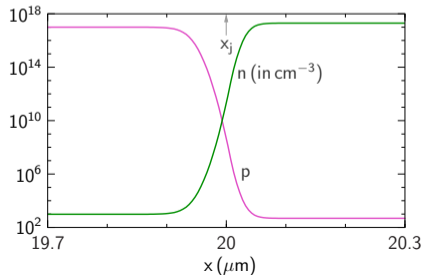
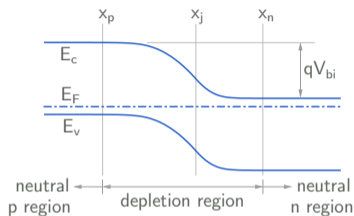
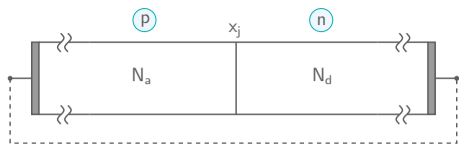
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* Qualitatively, we can see that the diffusion and drift currents will be in opposite directions:

Electrons:

$$J_n^{\text{diff}} : \leftarrow, \quad \mathcal{E} : \leftarrow, \quad J_n^{\text{drift}} : \rightarrow.$$



pn junction in equilibrium: current densities

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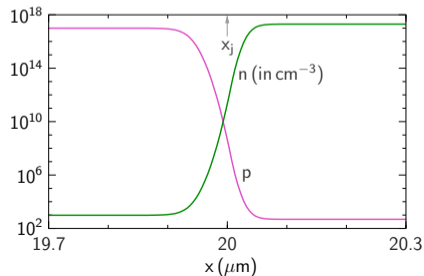
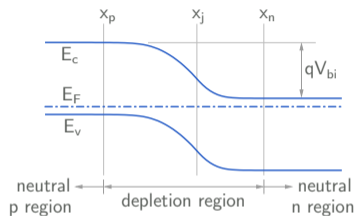
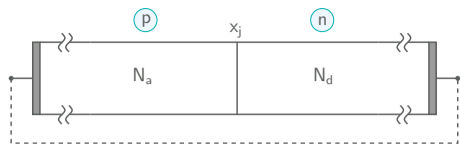
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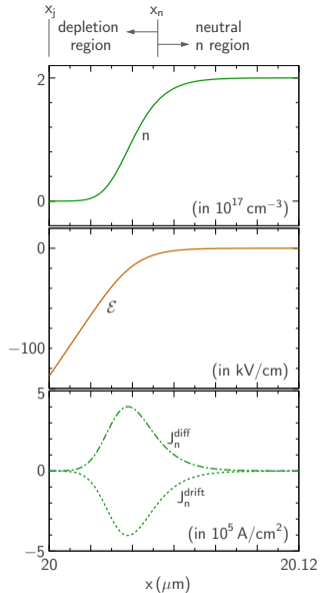
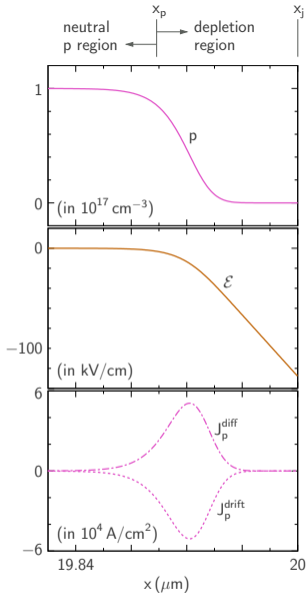
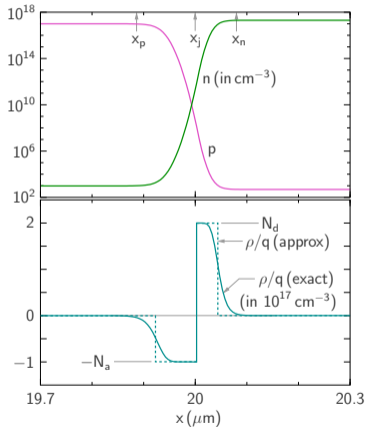
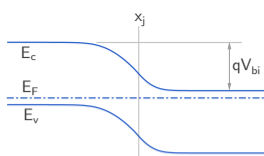
$$J_n^{\text{diff}} : \leftarrow, \quad \mathcal{E} : \leftarrow, \quad J_n^{\text{drift}} : \rightarrow.$$

Holes:

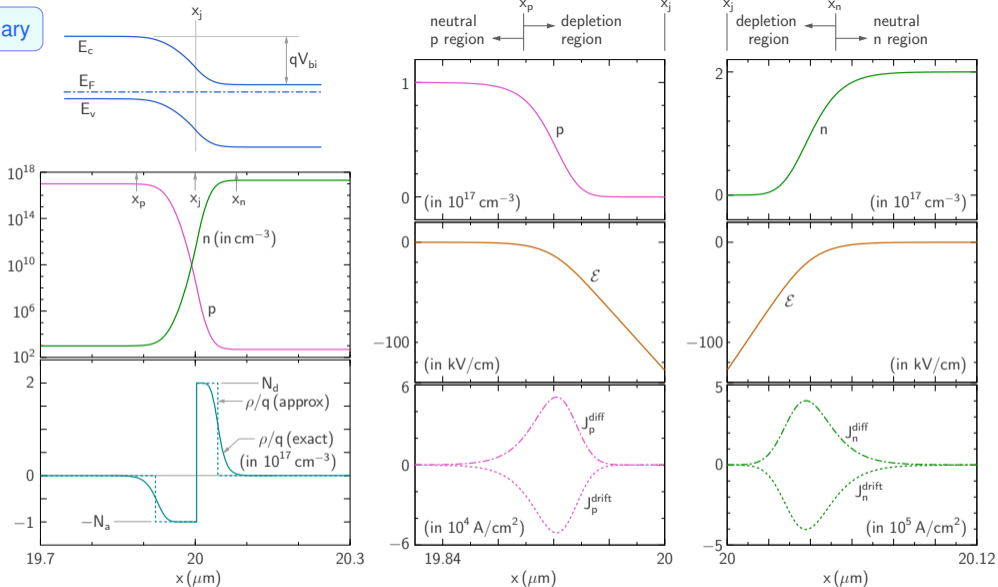
$$J_p^{\text{diff}} : \rightarrow, \quad \mathcal{E} : \leftarrow, \quad J_p^{\text{drift}} : \leftarrow.$$



Summary

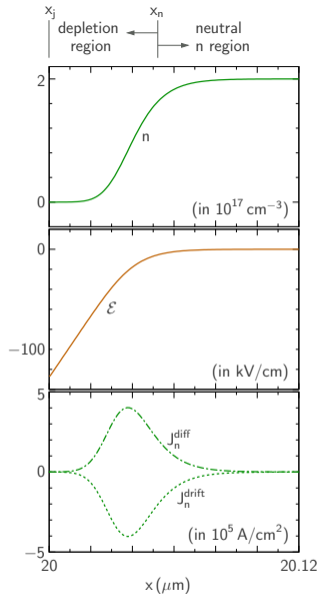
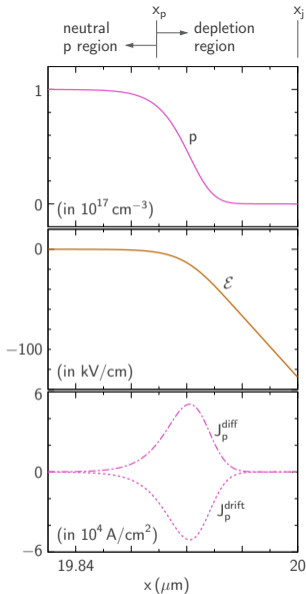
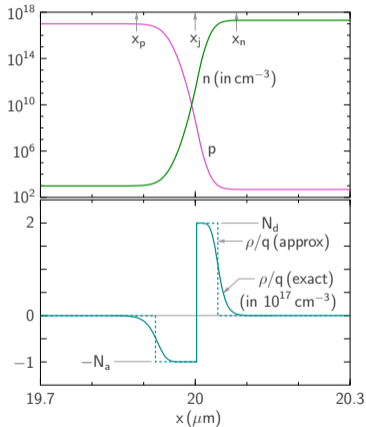
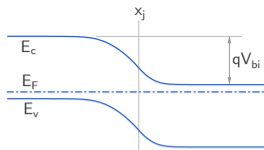


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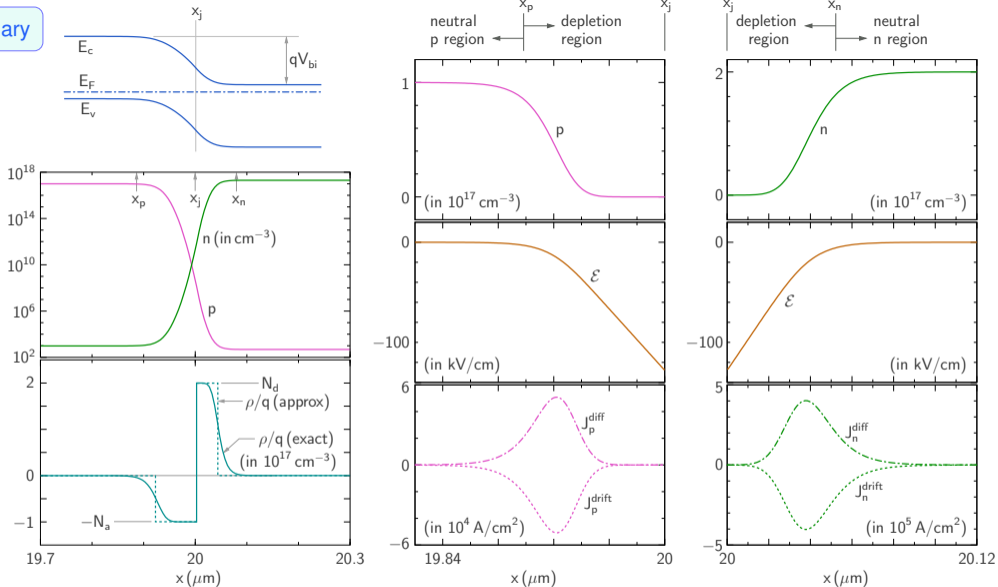


* There are three regions: p neutral region, n neutral region, and depletion region.

Summary

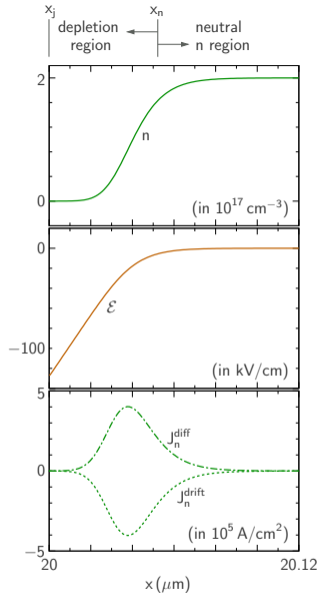
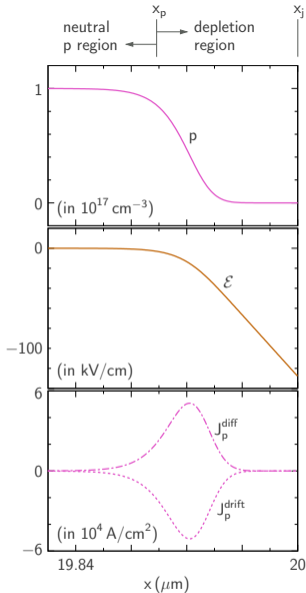
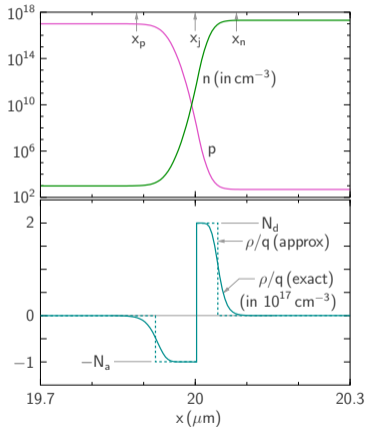
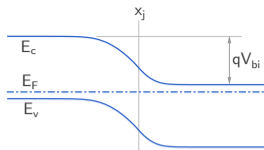


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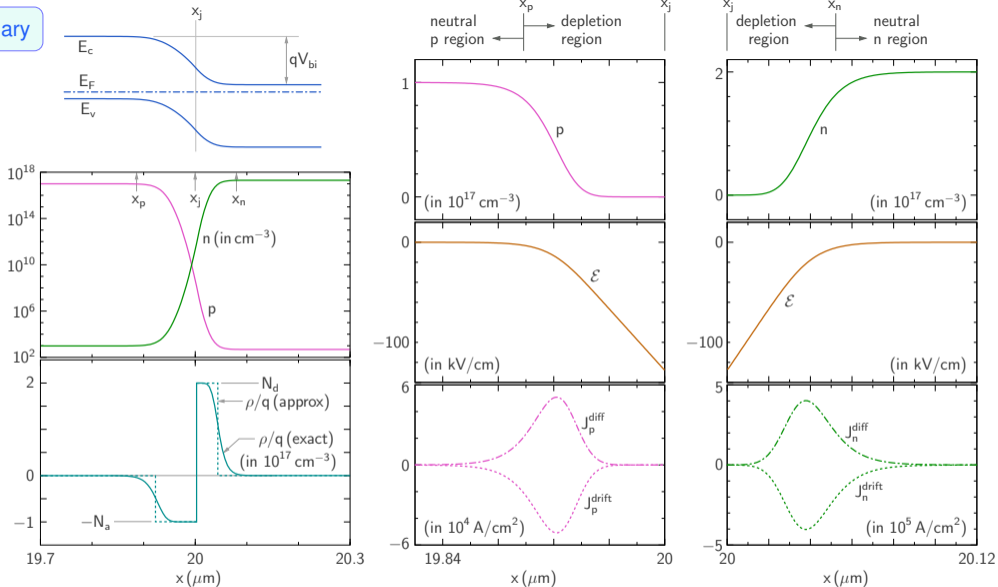


* The electric field is zero in the neutral regions and maximum (in magnitude) at the junction.

Summary

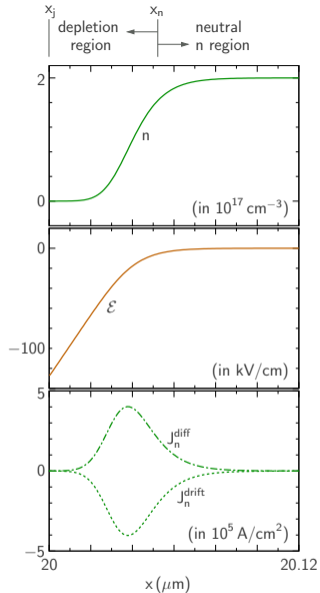
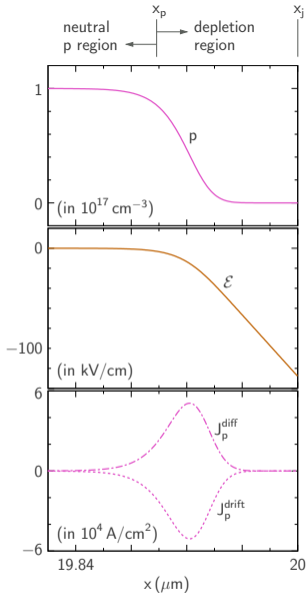
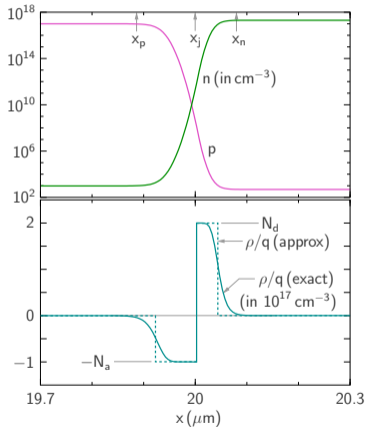
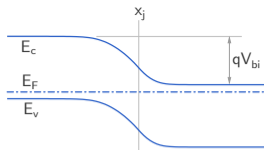


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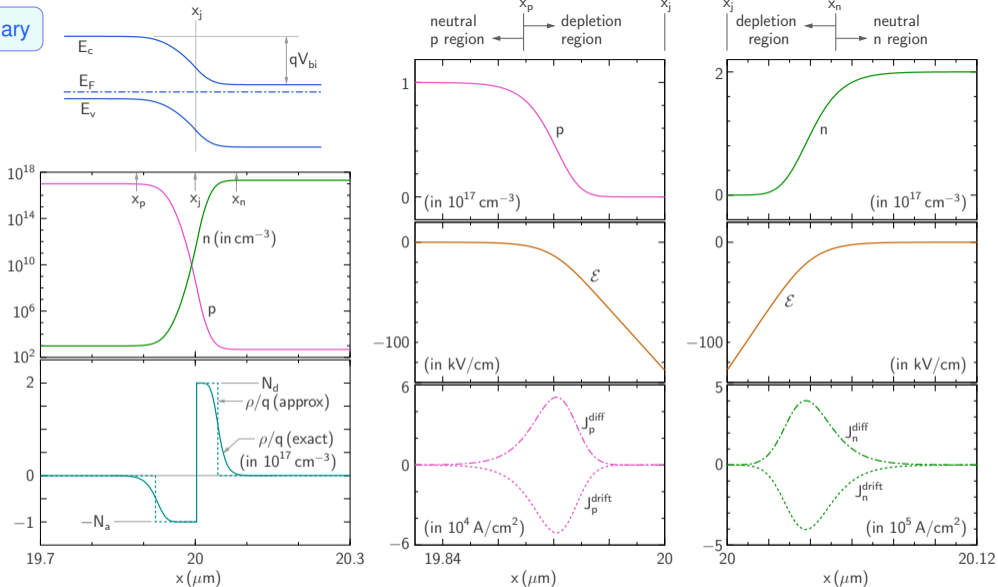


* There is a potential difference – the built-in voltage V_{bi} – between the neutral p and neutral n sides.

Summary



Summary



* J_n and J_p are individually zero because the drift and diffusion components cancel out.